

LECTURES

LESSON VI

Chapter 6. Special Functions

6.0 Introduction There is nothing particularly special about a special function, except that some person in authority or textbook writer (not the same thing!) has decided to bestow the moniker. Special functions are sometimes called higher transcendental functions (higher than what?) or functions of mathematical physics (but they occur in other fields also) or functions that satisfy certain frequently occurring second-order differential equations (but not all special functions do). One might simply call them useful functions and let it go at that; it is surely only a matter of taste which functions we have chosen to include in this chapter. Good commercially available program libraries, such as NAG or IMSL, contain routines for a number of special functions. These routines are intended for users who will have no idea what goes on inside them. Such state of the art black boxes are often very messy things, full of branches to completely different methods depending on the value of the calling arguments. Black boxes have, or should have, careful control of accuracy, to some stated uniform precision in all regimes. We will not be quite so fastidious in our examples, in part because we want to illustrate techniques

from Chapter 5, and in part because we want you to understand what goes on in the routines presented. Some of our routines have an accuracy parameter that can be made as small as desired, while others (especially those involving polynomial fits) give only a certain accuracy, one that we believe serviceable (typically six significant figures or more). We do not certify that the routines are perfect black boxes. We do hope that, if you ever encounter trouble in a routine, you will be able to diagnose and correct the problem on the basis of the information that we have given. In short, the special function routines of this chapter are meant to be used — we use them all the time but we also want you to be prepared to understand their inner workings.

CITED REFERENCES AND FURTHER READING: Abramowitz, M., and Stegun, I.A. 1964, Handbook of Mathematical Functions, Applied Mathematics Series, Volume 55 (Washington: National Bureau of Standards; reprinted 1968 by Dover Publications, New York) [full of useful numerical approximations to a great variety of functions]. IMSL Sfun/Library Users Manual (IMSL Inc., 2500 CityWest Boulevard, Houston TX 77042). NAG Fortran Library (Numerical Algorithms Group, 256 Banbury Road, Oxford OX27DE, U.K.), Chapter S. 205

206 Chapter 6. Special Functions Hart, J.F., et al. 1968, Computer Approximations (New York: Wiley). Hastings, C. 1955, Approximations for Digital Computers (Princeton: Princeton University Press). Luke, Y.L. 1975, Mathematical Functions and Their Approximations (New York: Academic Press).

6.1 Gamma Function, Beta Function, Factorials, Binomial Coefficients The gamma function is defined by the integral

$$\int_0^{\infty} t^{z-1} e^{-t} dt \quad (6.1.1) \quad \Gamma(z) = \frac{1}{z} \int_0^{\infty} t^{z-1} e^{-t} dt$$

When the argument z is an integer, the gamma function is just the familiar factorial function, but offset by one, $n! = \Gamma(n+1)$ (6.1.2) The gamma function

satisfies the recurrence relation $G(z+1) = zG(z)$ (6.1.3) If the function is known for arguments $z \geq 1$ or, more generally, in the half complex plane $\text{Re}(z) \geq 1$ it can be obtained for $z \geq 1$ or $\text{Re}(z) \geq 1$ by the reflection formula $G(1-z) = \frac{1}{\sin(\pi z)} G(z)$ (6.1.4) Notice that $G(z)$ has a pole at $z=0$, and at all negative integer values of z . There are a variety of methods in use for calculating the function $G(z)$ numerically, but none is quite as neat as the approximation derived by Lanczos [1]. This scheme is entirely specific to the gamma function, seemingly plucked from thin air. We will not attempt to derive the approximation, but only state the resulting formula: For certain integer choices of N , and for certain coefficients c_1, c_2, \dots, c_N , the gamma function is given by $G(z+1) = (z+1)^{-z} e^{-z} \left(\frac{2\pi}{z} \right)^{1/2} \left(1 + \frac{c_1}{z} + \frac{c_2}{z^2} + \dots + \frac{c_N}{z^N} \right)$ (6.1.5) You can see that this is a sort of take-off on Stirling's approximation, but with a series of corrections that take into account the first few poles in the left complex plane. The constant c_0 is very nearly equal to 1. The error term is parametrized by ϵ . For $N=5$, $\epsilon=10^{-10}$, and a certain set of c 's, the error is smaller than 10^{-10} . Impressed? If not, then perhaps you will be impressed by the fact that (with these

206 Chapter 6. Special Functions Hart, J.F., et al. 1968, Computer Approximations (New York: Wiley). Hastings, C. 1955, Approximations for Digital Computers (Princeton: Princeton University Press). Luke, Y.L. 1975, Mathematical Functions and Their Approximations (New York: Academic Press). 6.1 Gamma Function, Beta Function, Factorials, Binomial Coefficients The gamma function is defined by the integral $G(z) = \int_0^\infty t^{z-1} e^{-t} dt$ (6.1.1) When the argument z is an integer, the gamma function is just the familiar factorial function, but offset by one, $n! = G(n+1)$ (6.1.2) The gamma function satisfies the recurrence relation $G(z+1) = zG(z)$ (6.1.3) If the

function is known for arguments $z \geq 1$ or, more generally, in the half complex plane $\text{Re}(z) \geq 1$ it can be obtained for $z \geq 1$ or $\text{Re}(z) \geq 1$ by the reflection formula $\Gamma(1 - z) = \frac{1}{\Gamma(z) \sin(\pi z)}$ (6.1.4) Notice that $\Gamma(z)$ has a pole at $z = 0$, and at all negative integer values of z . There are a variety of methods in use for calculating the function $\Gamma(z)$ numerically, but none is quite as neat as the approximation derived by Lanczos [1]. This scheme is entirely specific to the gamma function, seemingly plucked from thin air. We will not attempt to derive the approximation, but only state the resulting formula: For certain integer choices of n and N , and for certain coefficients c_1, c_2, \dots, c_N , the gamma function is given by $\Gamma(z + 1) = \frac{e^{-z}}{z} \left(z + \frac{1}{2} + \frac{c_1}{z + 1} + \frac{c_2}{z + 2} + \dots + \frac{c_N}{z + N} \right)^z$ (6.1.5) You can see that this is a sort of take-off on Stirling's approximation, but with a series of corrections that take into account the first few poles in the left complex plane. The constant c_0 is very nearly equal to 1. The error term is parametrized by ϵ . For $n = 5$, $N = 6$, and a certain set of c 's, the error is smaller than 10^{-10} . Impressed? If not, then perhaps you will be impressed by the fact that (with these

6.1 Gamma, Beta, and Related Functions 207 same parameters) the formula (6.1.5) and bound on apply for the complex gamma function, everywhere in the half complex plane $\text{Re}(z) \geq 0$. It is better to implement $\ln \Gamma(x)$ than $\Gamma(x)$, since the latter will overflow many computers floating-point representation at quite modest values of x . Often the gamma function is used in calculations where the large values of $\Gamma(x)$ are divided by other large numbers, with the result being a perfectly ordinary value. Such operations would normally be coded as subtraction of logarithms. With (6.1.5) in hand, we can compute the logarithm of the gamma function with two calls to a logarithm and 25 or so arithmetic operations. This makes it not much

more difficult than other built-in functions that we take for granted, such as $\sin x$ or e^x :
FUNCTION gammln(xx) **REAL** gammln,xx Returns the value $\ln[\Gamma(xx)]$ for $xx > 0$.
INTEGER j **DOUBLE PRECISION** ser,stp,tmp,x,y,cof(6) Internal arithmetic will be done in double precision, a nicety that you can omit if five-figure accuracy is good enough. **SAVE** cof,stp
DATA cof,stp/76.18009172947146d0,-86.50532032941677d0, * 24.01409824083091d0,-1.231739572450155d0,.1208650973866179d0-
 2, * -5.395239384953d-5,2.5066282746310005d0/ x=xx y=x
 tmp=x+5.5d0 tmp=(x+0.5d0)*log(tmp)-tmp ser=1.000000000190015d0
 do 1 1 j=1,6 y=y+1.d0 ser=ser+cof(j)/y enddo 1 1 gammln=tmp+log(stp)
 return **END** How shall we write a routine for the factorial function $n!$? Generally the factorial function will be called for small integer values (for large values it will overflow anyway!), and in most applications the same integer value will be called for many times. It is a prodigal waste of computer time to call $\exp(\text{gammln}(n+1.0))$ for each required factorial. Better to go back to basics, holding gammln in reserve for unlikely calls:
FUNCTION factrl(n) **INTEGER** n **REAL** factrl C USES gammln Returns the value $n!$ as a floating-point number. **INTEGER** j,ntop **REAL** a(33),gammln Table to be filled in only as required. **SAVE** ntop,a **DATA** ntop,a(1)/0,1./ Table initialized with $0!$ only. if (n.lt.0) then pause negative factorial in factrl else if (n.le.ntop) then Already in table. factrl=a(n+1) else if (n.le.32) then Fill in table up to desired value. do 1 1 j=ntop+1,n

208 Chapter 6. Special Functions $a(j+1)=j*a(j)$ enddo 1 1 ntop=n factrl=a(n+1) else Larger value than size of table is required. Actually, this big a value is going to overflow on many computers, but no harm in trying. factrl= $\exp(\text{gammln}(n+1.))$ endif return **END** A useful point is that factrl will be exact for the smaller values of n , since floating-point multiplies on small

integers are exact on all computers. This exactness will not hold if we turn to the logarithm of the factorials. For binomial coefficients, however, we must do exactly this, since the individual factorials in a binomial coefficient will grow long before the coefficient itself will. The binomial coefficient is defined by $\binom{n}{k} = \frac{n!}{k!(n-k)!}$ $0 \leq k \leq n$ (6.1.6)

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FUNCTION bico(n,k) INTEGER k,n REAL bico C USES factln
Returns the binomial coefficient  $\binom{n}{k}$  as a floating-point number.
REAL factln bico=nint(exp(factln(n)-factln(k)-factln(n-k)))
return The nearest-integer function cleans up roundoff error for smaller values of n and k.
END which uses
FUNCTION factln(n) INTEGER n REAL factln C USES gammln
Returns ln(n!).
REAL a(100),gammln SAVE a DATA a/100*-1./
Initialize the table to negative values.
if (n.lt.0) pause negative factorial in factln
if (n.le.99) then In range of the table.
if (a(n+1).lt.0.) a(n+1)=gammln(n+1.)
If not already in the table, put it in.
factln=a(n+1) else factln=gammln(n+1.)
Out of range of the table. endif return
END

```

209 6.2 Incomplete Gamma Function If your problem requires a series of related binomial coefficients, a good idea is to use recurrence relations, for example $\binom{n}{k} = \frac{n}{k} \binom{n-1}{k-1}$ $\binom{n}{k} = \frac{n-k+1}{k} \binom{n}{k-1}$ (6.1.7) $\binom{n}{k} \binom{n-k}{k+1} = \binom{n-k}{k} \binom{n}{k+1}$

Finally, turning away from the combinatorial functions with integer valued arguments, we come to the beta function, $\int_0^1 (1-t)^{w-1} t^{z-1} dt$ (6.1.8) $B(z, w) = B(w, z) = \frac{\Gamma(z)\Gamma(w)}{\Gamma(z+w)}$ (6.1.9) $G(z+w)$ hence

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FUNCTION beta(z,w) REAL beta,w,z C USES gammln
Returns the value of the beta function B(z, w).
REAL gammln beta=exp(gammln(z)+gammln(w)-gammln(z+w))
return
END

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CITED REFERENCES AND FURTHER READING: Abramowitz, M., and Stegun, I.A. 1964, Handbook of Mathematical Functions, Applied Mathematics Series, Vol-

ume 55 (Washington: National Bureau of Standards; reprinted 1968 by Dover Publications, New York), Chapter 6. Lanczos, C. 1964, SIAM Journal on Numerical Analysis, ser. B, vol. 1, pp. 8696. [1] 6.2 Incomplete Gamma Function, Error Function, Chi-Square Probability Function, Cumulative Poisson Function The incomplete gamma function is defined by $x P(a, x) = (a, x) e^{-x} x^{a-1} dt (a; 0)$ (6.2.1) $G(a) = 1 G(a) 0$

209 6.2 Incomplete Gamma Function If your problem requires a series of related binomial coefficients, a good idea is to use recurrence relations, for example $n k = n n + 1 k = n + 1 n - k + 1 k + n k - 1$ (6.1.7) $n k n k + 1 = n - k k + 1$ Finally, turning away from the combinatorial functions with integer valued arguments, we come to the beta function, $t z^{-1} (1 - t)^{w-1} dt$ (6.1.8) $B(z, w) = B(w, z) = \frac{1}{0}$ which is related to the gamma function by $B(z, w) = \frac{G(z)G(w)}{G(z+w)}$ (6.1.9) $G(z + w)$ hence FUNCTION beta(z,w) REAL beta,w,z C USES gammln Returns the value of the beta function $B(z, w)$. REAL gammln beta= $\exp(\text{gammln}(z) + \text{gammln}(w) - \text{gammln}(z+w))$ return END CITED REFERENCES AND FURTHER READING: Abramowitz, M., and Stegun, I.A. 1964, Handbook of Mathematical Functions, Applied Mathematics Series, Volume 55 (Washington: National Bureau of Standards; reprinted 1968 by Dover Publications, New York), Chapter 6. Lanczos, C. 1964, SIAM Journal on Numerical Analysis, ser. B, vol. 1, pp. 8696. [1] 6.2 Incomplete Gamma Function, Error Function, Chi-Square Probability Function, Cumulative Poisson Function The incomplete gamma function is defined by $x P(a, x) = (a, x) e^{-x} x^{a-1} dt (a; 0)$ (6.2.1) $G(a) = 1 G(a) 0$

210 Chapter 6. Special Functions 1.0 0.5 1.0 $a = 3.0$ $a = 10$.2 .4 .6 .8 0 02 4 6 8 10 12 14 x Figure 6.2.1. The incomplete gamma function $P(a, x)$ for four values of a . It has the limiting values $P(a, 0) = 0$ and $P(a, \infty) = 1$ (6.2.2) The incomplete

gamma function $P(a, x)$ is monotonic and (for a greater than one or so) rises from near-zero to near-unity in a range of x centered on about $a - 1$, and of width about \sqrt{a} (see Figure 6.2.1). The complement of $P(a, x)$ is also confusingly called an incomplete gamma function, $Q(a, x) = 1 - P(a, x) = \frac{1}{\Gamma(a)} \int_x^\infty t^{a-1} e^{-t} dt$ (6.2.3) $\Gamma(a) = \int_0^\infty t^{a-1} e^{-t} dt$ (6.2.4) It has the limiting values $Q(a, 0) = 1$ and $Q(a, \infty) = 0$ (6.2.4) The notations $P(a, x)$, $Q(a, x)$, and $\Gamma(a, x)$ are standard; the notation $Q(a, x)$ is specific to this book. There is a series development for $Q(a, x)$ as follows: $Q(a, x) = e^{-x} x^a \sum_{n=0}^{\infty} \frac{(-1)^n x^n}{\Gamma(a+1+n)}$ (6.2.5) One does not actually need to compute a new $\Gamma(a+1+n)$ for each n ; one rather uses equation (6.1.3) and the previous coefficient.

211 6.2 Incomplete Gamma Function A continued fraction development for $Q(a, x)$ is $Q(a, x) = e^{-x} x^a \frac{1}{x+1} \frac{1-a}{x+2} \frac{1-a}{x+3} \dots$ (6.2.6) It is computationally better to use the even part of (6.2.6), which converges twice as fast (see 5.2): $Q(a, x) = e^{-x} x^a \frac{1-a}{x+1} \frac{1-a}{x+3} \frac{1-a}{x+5} \dots$ (6.2.7) It turns out that (6.2.5) converges rapidly for x less than about $a+1$, while (6.2.6) or (6.2.7) converges rapidly for x greater than about $a+1$. In these respective regimes each requires at most a few times \sqrt{a} terms to converge, and this many only near $x = a$, where the incomplete gamma functions are varying most rapidly. Thus (6.2.5) and (6.2.7) together allow evaluation of the function for all positive a and x . An extra dividend is that we never need compute a function value near zero by subtracting two nearly equal numbers. The higher-level functions that return $P(a, x)$ and $Q(a, x)$ are `FUNCTION gammp(a,x) REAL a,gammp,x C USES gcf,gser Returns the incomplete gamma function P(a, x). REAL gammcf,gamser,gl n if(x.lt.0..or.a.le.0.)pause bad arguments in gammp if(x.lt.a+1.)then Use the series representation. call gser(gamser,a,x,gl n) gammp=gamser else Use`


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the continued fraction representation call gcf(gammcf,a,x,gln)
gamp=1.-gammcf and take its complement. endif return
END FUNCTION gammq(a,x) REAL a,gammq,x C USES
gcf,gser Returns the incomplete gamma function  $Q(a, x) = 1 - P(a, x)$ . REAL gammcf,gamser,gln if(x.lt.0..or.a.le.0.)pause
bad arguments in gammq if(x.lt.a+1.)then Use the series representation call gser(gamser,a,x,gln) gammq=1.-gamser and take
its complement. else Use the continued fraction representation. call gcf(gammcf,a,x,gln) gammq=gammcf endif return END

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212 Chapter 6. Special Functions The argument gln is returned by both the series and continued fraction procedures containing the value $\ln \Gamma(a)$; the reason for this is so that it is available to you if you want to modify the above two procedures to give (a, x) and $\Gamma(a, x)$, in addition to $P(a, x)$ and $Q(a, x)$ (cf. equations 6.2.1 and 6.2.3). The procedures gser and gcf which implement (6.2.5) and (6.2.7) are SUBROUTINE gser(gamser,a,x,gln) INTEGER ITMAX REAL a,gamser,gln,x,EPS PARAMETER (ITMAX=100,EPS=3.e-7) C USES gammln Returns the incomplete gamma function $P(a, x)$ evaluated by its series representation as gamser. Also returns $\ln \Gamma(a)$ as gln. INTEGER n REAL ap,del,sum,gammln gln=gammln(a) if(x.le.0.)then if(x.lt.0.)pause x ≤ 0 in gser gamser=0. return endif ap=a sum=1./a del=sum do 1 1 n=1,ITMAX ap=ap+1. del=del*x/ap sum=sum+del if(abs(del).lt.abs(s 1 enddo 1 1 pause a too large, ITMAX too small in gser 1 gamser=sum*exp(-x+a*log(x)-gln) return END SUBROUTINE gcf(gammcf,a,x,gln) INTEGER ITMAX REAL a,gammcf,gln,x,EPS PARAMETER (ITMAX=100,EPS=3.e-7,FPMIN=1.e-30) C USES gammln Returns the incomplete gamma function $Q(a, x)$ evaluated by its continued fraction representation as gammcf. Also returns $\ln \Gamma(a)$ as gln. Parameters: ITMAX is the maximum allowed number of iterations; EPS is the

relative accuracy; FPMIN is a number near the smallest representable floating-point number. INTEGER i REAL an,b,c,d,del,h,gammln gln=gammln(a) b=x+1.-a Set up for evaluating continued fraction by modified Lentz's method (5.2) with b0 =0. c=1./FPMIN d=1./b h=d do 1 1 i=1,ITMAX Iterate to convergence. an=-i*(i-a) b=b+2. d=an*d+b if(abs(d).lt.FPMIN)d=FPMIN c=b+an/c if(abs(c).lt.FPMIN)c=FPMIN d=1./d del=d*c

213 6.2 Incomplete Gamma Function h=h*del if(abs(del-1.).lt.EPS)goto 1 enddo 1 1 pause a too large, ITMAX too small in gcf 1 gammcf=exp(-x+a*log(x)-gln)*h Put factors in front. return END Error Function The error function and complementary error function are special cases of the incomplete gamma function, and are obtained moderately efficiently by the above procedures. Their definitions are $\text{erf}(x) = \frac{2}{\sqrt{\pi}} \int_0^x e^{-t^2} dt$ (6.2.8) and $\text{erfc}(x) = 1 - \text{erf}(x) = \frac{2}{\sqrt{\pi}} \int_x^\infty e^{-t^2} dt$ (6.2.9) x The functions have the following limiting values and symmetries: $\text{erf}(0) = 0$ $\text{erf}(\infty) = 1$ $\text{erf}(-x) = -\text{erf}(x)$ (6.2.10) $\text{erfc}(0) = 1$ $\text{erfc}(\infty) = 0$ $\text{erfc}(-x) = 2 - \text{erfc}(x)$ (6.2.11) They are related to the incomplete gamma functions by $\text{erf}(x) = P(1/2, x^2)$ ($x = 0$) (6.2.12) and $\text{erfc}(x) = Q(1/2, x^2)$ ($x = 0$) (6.2.13) Hence we have FUNCTION erf(x) REAL erf,x C USES gammp Returns the error function erf(x). REAL gammp if(x.lt.0.)then erf=-gammp(.5,x**2) else erf=gammp(.5,x**2) endif return END

215 6.3 Exponential Integrals Chi-Square Probability Function $P(\chi^2, \nu)$ is defined as the probability that the observed chi-square for a correct model should be less than a value χ^2 . (We will discuss the use of this function in Chapter 15.) Its complement $Q(\chi^2, \nu)$ is the probability that the observed chi-square will exceed the value χ^2 by chance even for a correct model. In both cases ν is an integer, the number of degrees of freedom. The functions have the limiting values $P(0, \nu) = 0$ $P(\infty, \nu) = 1$

$(8 -) = 1$ (6.2.16) $Q(0 -) = 1$ $Q(8 -) = 0$ (6.2.17) and the following relation to the incomplete gamma functions, $P(2 -) = P^2 = \text{gammp } 2$ (6.2.18) $2, 2, 2$ $Q(2 -) = Q^2 = \text{gammq } 2$ (6.2.19) $2, 2, 2$

CITED REFERENCES AND FURTHER READING: Abramowitz, M., and Stegun, I.A. 1964, Handbook of Mathematical Functions, Applied Mathematics Series, Volume 55 (Washington: National Bureau of Standards; reprinted 1968 by Dover Publications, New York), Chapters 6, 7, and 26. Pearson, K. (ed.) 1951, Tables of the Incomplete Gamma Function (Cambridge: Cambridge University Press).

6.3 Exponential Integrals The standard definition of the exponential integral is $e^{-x} \int_0^x \frac{1}{t} dt$, $x \geq 0, n=0, 1, \dots$ (6.3.1) The function defined by the principal value of the integral $\int_0^x \frac{e^{-t}}{t} dt$, $x \geq 0$ (6.3.2) $\int_0^x \frac{e^{-t}}{t} dt = x^{-x} - \gamma$ is also called an exponential integral. Note that $Ei(-x)$ is related to $-E_1(x)$ by analytic continuation. The function $E_n(x)$ is a special case of the incomplete gamma function $E_n(x) = x^{n-1} \Gamma(1-n, x)$ (6.3.3)

216 Chapter 6. Special Functions We can therefore use a similar strategy for evaluating it. The continued fraction just equation (6.2.6) rewritten converges for all $x \geq 0$: $E_n(x) = e^{-x} \frac{1}{1 + \frac{x}{1 + \frac{x}{1 + \frac{x}{\dots}}}}$ (6.3.4) We use it in its more rapidly converging even form, $1 - \frac{x}{2(n+1)} \frac{E_n(x)}{E_{n+1}(x)} = e^{-x} \frac{1}{1 + \frac{x}{n} - \frac{x}{n+2} + \frac{x}{n+4} - \dots}$ (6.3.5) The continued fraction only really converges fast enough to be useful for $x \geq 1$. For $0 \leq x < 1$, we can use the series representation $\sum_{m=0}^{\infty} \frac{(-x)^m}{m!} E_n(x) = \frac{(-x)^{n-1}}{(n-1)!} [-\ln x + \psi(n)] - \frac{(m-n+1)m!}{(n-1)!} \dots$ (6.3.6) $m=0$ $m=n-1$ The quantity $\psi(n)$ here is the digamma function, given for integer arguments by $\psi(1) = -\gamma$, $\psi(n) = -\gamma + \sum_{k=1}^{n-1} \frac{1}{k}$ (6.3.7) $m=1$ where $\gamma = 0.5772156649\dots$ is Euler's constant. We evaluate the expression (6.3.6) in order of ascending powers of x : $E_n(x) = -\frac{1}{(n-1)!} (1-n) - x \frac{(-1)^{(n-2)!}}{(2-n)!} + x^2 \frac{(3-n)(1-2)}{(3-n)!} + \dots + (-x)^{n-2} \frac{(-1)^{(n-1)!}}{(n-1)!} + (-x)^{n-1} \frac{(-1)^{(n-1)!}}{(n-1)!} [-\ln x + \psi(n)] - \frac{(-x)^n}{n!} + \dots$

(-x)ⁿ + 1 (6.3.8) The first square bracket is omitted when n = 1. This method of evaluation has the advantage that for large n the series converges before reaching the term containing (n). Accordingly, one needs an algorithm for evaluating (n) only for small n, n ≤ 20–40. We use equation (6.3.7), although a table look-up would improve efficiency slightly. Amos [1] presents a careful discussion of the truncation error in evaluating equation (6.3.8), and gives a fairly elaborate termination criterion. We have found that simply stopping when the last term added is smaller than the required tolerance works about as well. Two special cases have to be handled separately: E₀(x) = e^{-x} x (6.3.9) E_n(0) = 1/n - 1, n ≥ 1

217 6.3 Exponential Integrals The routine expint allows fast evaluation of E_n(x) to any accuracy EPS within the reach of your machines word length for floating-point numbers. The only modification required for increased accuracy is to supply Eulers constant with enough significant digits. Wrench [2] can provide you with the first 328 digits if necessary! FUNCTION expint(n,x) INTEGER n,MAXIT REAL expint,x,EPS,FPMIN,EULER PARAMETER (MAXIT=100,EPS=1.e-7,FPMIN=1.e-30,EULER=.5772156649) Evaluates the exponential integral E_n(x). Parameters: MAXIT is the maximum allowed number of iterations; EPS is the desired relative error, not smaller than the machine precision; FPMIN is a number near the smallest representable floating-point number; EULER is Eulers constant. INTEGER i,ii,nm1 REAL a,b,c,d,del,fact,h,psi nm1=n-1 if(n.lt.0.or.x.lt.0..or.(x.eq.0..and.(n.eq.0.or pause bad arguments in expint else if(n.eq.0)then Special case. expint=exp(-x)/x else if(x.eq.0.)then Another special case. expint=1./nm1 else if(x.gt.1.)then Lentzs algorithm (5.2). b=x+n c=1./FPMIN d=1./b h=d do 1 1 i=1,MAXIT a=-i*(nm1+i) b=b+2. d=1./(a*d+b) Denominators cannot

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be zero. c=b+a/c del=c*d h=h*del if(abs(del-1.).lt.EPS)then
expint=h*exp(-x) return endif enddo 1 1 pause continued frac-
tion failed in expint else Evaluate series. if(nm1.ne.0)then Set-
firstterm. expint=1./nm1 else expint=-log(x)-EULER endif
fact=1. do 1 3 i=1,MAXIT fact=-fact*x/i if(i.ne.nm1)then
del=-fact/(i-nm1) else psi=-EULER Compute (n). do 1 2
ii=1,nm1 psi=psi+1./ii enddo 1 2 del=fact*(-log(x)+psi) en-
dif expint=expint+del if(abs(del).lt.abs(expint)*EPS) return
enddo 1 3

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pause series failed in
expint endif return END

A good algorithm for evaluating Ei is to use the power series for small x and the asymptotic series for large x . The power series is $Ei(x) = -\ln x + x/2! + x^2/3! + \dots$ (6.3.10) where γ is Euler's constant. The asymptotic expansion is $Ei(x) \sim \frac{e^x}{x} \left(1 + \frac{\gamma}{x} + \frac{1}{2x^2} + \dots \right)$ (6.3.11)

The lower limit for the use of the asymptotic expansion is approximately $-\ln \text{EPS}$, where EPS is the required relative error.

```

FUNCTION ei(x) INTEGER MAXIT
REAL ei,x,EPS,EULER,FPMIN PARAMETER (EPS=6.e-8,EULER=.57721566,MAXIT=100,FPMIN=1.e-30)
Computes the exponential integral Ei(x) for x>0. Parameters: EPS is the relative error, or absolute error near the zero of Ei at x=0.3725; EULER is Euler's constant; MAXIT is the maximum number of iterations allowed; FPMIN is a number near the smallest representable floating-point number.
INTEGER k REAL fact,prev,sum,term if(x.le.0.) pause bad argument in ei if(x.lt.FPMIN)then Special case: avoid failure of convergence test because of underflow. ei=log(x)+EULER else if(x.le.-log(EPS))then Use power series. sum=0. fact=1. do 1 1 k=1,MAXIT fact=fact*x/k term=fact/k sum=sum+term if(term.lt.EPS*sum)goto 1 enddo 1 1 pause series failed in ei 1 ei=sum+log(x)+EULER else Use asymptotic series. sum=0.

```

```

Start with second term. term=1. do 1 2 k=1,MAXIT
prev=term term=term*k/x if(term.lt.EPS)goto 2 Since final
sum is greater than one, term itself approximates the rel-
ative error. if(term.lt.prev)then sum=sum+term Still con-
vergi ng: add new term. else sum=sum-prev Diverging: sub-
tract previous term and exit. goto 2 endif enddo 1 2 2
ei=exp(x)*(1.+sum)/x endif

```

219 6.4 Incomplete Beta Function, Students Distribution, F-Distribution, Cumulative Binomial Distribution return END

CITED REFERENCES AND FURTHER READING: Stegun, I.A., and Zucker, R. 1974, Journal of Research of the National Bureau of Standards, vol. 78B, pp. 199216; 1976, op. cit., vol. 80B, pp. 291311. Amos D.E. 1980, ACM Transactions on Mathematical Software, vol. 6, pp. 365377 [1]; also vol. 6, pp. 420428. Abramowitz, M., and Stegun, I.A. 1964, Handbook of Mathematical Functions, Applied Mathematics Series, Volume 55 (Washington: National Bureau of Standards; reprinted 1968 by Dover Publications, New York), Chapter 5. Wrench J.W. 1952, Mathematical Tables and Other Aids to Computation, vol. 6, p. 255. [2]

6.4 Incomplete Beta Function, Students Distribution, F-Distribution, Cumulative Binomial Distribution The incomplete beta function is defined by

$$I_x(a, b) = \frac{B_x(a, b)}{B(a, b)} = \int_0^x t^{a-1} (1-t)^{b-1} dt \quad (a, b > 0) \quad (6.4.1)$$

$$B(a, b) = \int_0^1 t^{a-1} (1-t)^{b-1} dt \quad (6.4.2)$$

It has the limiting values $I_0(a, b) = 0$ $I_1(a, b) = 1$ (6.4.2) and the symmetry relation $I_x(a, b) = 1 - I_{1-x}(b, a)$ (6.4.3) If a and b are both rather greater than one, then $I_x(a, b)$ rises from near-zero to near-unity quite sharply at about $x = a/(a + b)$. Figure 6.4.1 plots the function for several pairs (a, b) . The incomplete beta function has a series expansion

$$I_x(a, b) = \frac{x^a (1-x)^{b-a}}{a B(a, b)} \sum_{n=0}^{\infty} \frac{B(a+n, b-n)}{B(a, b)} x^n \quad (6.4.4)$$

but this does not prove to be very useful in its numerical evaluation. (Note, however, that the beta functions

in the coefficients can be evaluated for each value of n with just the previous value and a few multiplies, using equations 6.1.9 and 6.1.3.) The continued fraction representation proves to be much more useful, $I_x(a, b) = \frac{x^a (1-x)^b}{B(a, b)} \frac{1 + \frac{1-x}{1+x} \frac{1-x}{1+x} \dots}{1 + \frac{1-x}{1+x} \frac{1-x}{1+x} \dots}$ (6.4.5)

219 6.4 Incomplete Beta Function, Students Distribution, F-Distribution, Cumulative Binomial Distribution return END
 CITED REFERENCES AND FURTHER READING: Stegun, I.A., and Zucker, R. 1974, Journal of Research of the National Bureau of Standards, vol. 78B, pp. 199216; 1976, op. cit., vol. 80B, pp. 291311. Amos D.E. 1980, ACM Transactions on Mathematical Software, vol. 6, pp. 365377 [1]; also vol. 6, pp. 420428. Abramowitz, M., and Stegun, I.A. 1964, Handbook of Mathematical Functions, Applied Mathematics Series, Volume 55 (Washington: National Bureau of Standards; reprinted 1968 by Dover Publications, New York), Chapter 5. Wrench J.W. 1952, Mathematical Tables and Other Aids to Computation, vol. 6, p. 255. [2] 6.4 Incomplete Beta Function, Students Distribution, F-Distribution, Cumulative Binomial Distribution The incomplete beta function is defined by $I_x(a, b) = \frac{B_x(a, b)}{B(a, b)} = \frac{\int_0^x t^{a-1} (1-t)^{b-1} dt}{\int_0^1 t^{a-1} (1-t)^{b-1} dt}$ (6.4.1) $B(a, b) = \frac{\Gamma(a)\Gamma(b)}{\Gamma(a+b)}$ (6.4.2) It has the limiting values $I_0(a, b) = 0$ $I_1(a, b) = 1$ (6.4.2) and the symmetry relation $I_x(a, b) = 1 - I_{1-x}(b, a)$ (6.4.3) If a and b are both rather greater than one, then $I_x(a, b)$ rises from near-zero to near-unity quite sharply at about $x = a/(a + b)$. Figure 6.4.1 plots the function for several pairs (a, b) . The incomplete beta function has a series expansion $I_x(a, b) = \frac{x^a (1-x)^b}{B(a, b)} \frac{1 + \frac{1-x}{1+x} \frac{1-x}{1+x} \dots}{1 + \frac{1-x}{1+x} \frac{1-x}{1+x} \dots}$ (6.4.4) $n = 0$ but this does not prove to be very useful in its numerical evaluation. (Note, however, that the beta functions in the coefficients can be evaluated for each value of n with just the previous value and a few multiplies, using equations 6.1.9

and 6.1.3.) The continued fraction representation proves to be much more useful, $I_x(a, b) = \frac{x^a (1-x)^b}{B(a, b)}$ (6.4.5)

220 Chapter 6. Special Functions 1 (0.5,5.0) (8.0,10.0) (1.0,3.0) (0.5,0.5) .2 .4 .6 .8 (5.0,0.5) 0 0 .2 .4 .6 1 .8 x Figure 6.4.1. The incomplete beta function $I_x(a, b)$ for five different pairs of (a, b) . Notice that the pairs $(0.5, 5.0)$ and $(5.0, 0.5)$ are symmetrically related as indicated in equation (6.4.3). where $d_{2m+1} = - (a + m)(a + b + m)x (a + 2m)(a + 2m + 1)$ (6.4.6) $d_{2m} = m(b - m)x (a + 2m - 1)(a + 2m)$ This continued fraction converges rapidly for $x_i(a + 1)/(a + b + 2)$, taking in the worst case $O(\max(a, b))$ iterations. But for $x_i(a + 1)/(a + b + 2)$ we can just use the symmetry relation (6.4.3) to obtain an equivalent computation where the continued fraction will also converge rapidly. Hence we have

```

FUNCTION betai(a,b,x) REAL betai,a,b,x C USES betacf,gammln
Returns the incomplete beta function Ix(a, b).
REAL bt,betacf,gammln if(x.lt.0..or.x.gt.1.)pause bad argument x in betai if(x.eq.0..or.x.eq.1.)then bt=0. else
Factors in front of the continued fraction. bt=exp(gammln(a+b)-gammln(a)-gammln(b) * +a*log(x)+b*log(1.-x)) endif if(x.lt.(a+1.)/(a+
Use continued fraction directly.

```

221 6.4 Incomplete Beta Function, Students Distribution, F-Distribution, Cumulative Binomial Distribution $betai = bt * betacf(a, b, x) / a$ return else $betai = 1. - bt * betacf(b, a, 1. - x) / b$ Use continued fraction after making the symmetry transformation. return endif END which utilizes the continued fraction evaluation routine

```

FUNCTION betacf(a,b,x) INTEGER MAXIT REAL betacf,a,b,x,EPS,FPMIN PARAMETER (MAXIT=100,EPS=3.e-7,FPMIN=1.e-30)
Used by betai: Eval uates continued fraction for incomplete beta function by modified Lentzs method (5.2). INTEGER m,m2 REAL

```



```

aa,c,d,del,h,qab,qam,qap qab=a+b These qs will be used in
factors that occur in the coefficients (6.4.6). qap=a+1. qam=a-
1. c=1. First step of Lentz's method. d=1.-qab*x/qap
if(abs(d).lt.FPMIN)d=FPMIN d=1./d h=d do 1 1 m=1,MAXIT
m2=2*m aa=m*(b-m)*x/((qam+m2)*(a+m2)) d=1.+aa*d One
step (the even one) of the recurrence. if(abs(d).lt.FPMIN)
c=1.+aa/c if(abs(c).lt.FPMIN)c=FPMIN d=1./d h=h*d*c aa=-
(a+m)*(qab+m)*x/((a+m2)*(qap+m2)) d=1.+aa*d Next step
of the recurrence (the odd one). if(abs(d).lt.FPMIN)
c=1.+aa/c if(abs(c).lt.FPMIN)c=FPMIN d=1./d del=d*c h=h*del
if(abs(del-1.).lt.EPS)goto 1 Are we done? enddo 1 1 pause a
or b too big, or MAXIT too small in betacf 1 betacf=h return
END

```

Students Distribution Probability Function Students distribution, denoted $A(t; \nu)$, is useful in several statistical contexts, notably in the test of whether two observed distributions have the same mean. $A(t; \nu)$ is the probability, for degrees of freedom ν , that a certain statistic t (measuring the observed difference of means) would be smaller than the observed value if the means were in fact the same. (See Chapter 14 for further details.) Two means are

222 Chapter 6. Special Functions significantly different if, e.g., $A(t; \nu) \leq 0.99$. In other words, $1 - A(t; \nu)$ is the significance level at which the hypothesis that the means are equal is disproved. The mathematical definition of the function is

$$(6.4.7) \quad A(t; \nu) = \frac{1}{2} B\left(\frac{1-t}{2}, \frac{\nu}{2}\right) / B\left(\frac{1-t}{2}, \frac{\nu}{2}\right)$$

Limiting values are $A(0; \nu) = 0$, $A(\infty; \nu) = 1$ (6.4.8) $A(t; \nu)$ is related to the incomplete beta function $I_x(a, b)$ by

$$(6.4.9) \quad A(t; \nu) = 1 - I_{t^2/2}(1, \nu/2)$$

So, you can use (6.4.9) and the above routine `betainc` to evaluate the function. F-Distribution Probability Function This function occurs in the statistical test of whether two observed samples have the same

variance. A certain statistic F , essentially the ratio of the observed dispersion of the first sample to that of the second one, is calculated. (For further details, see Chapter 14.) The probability that F would be as large as it is if the first samples underlying distribution actually has smaller variance than the second is denoted $Q(F - 1, 2)$, where 1 and 2 are the number of degrees of freedom in the first and second samples, respectively. In other words, $Q(F - 1, 2)$ is the significance level at which the hypothesis 1 has smaller variance than 2 can be rejected. A small numerical value implies a very significant rejection, in turn implying high confidence in the hypothesis 1 has variance greater or equal to 2. $Q(F - 1, 2)$ has the limiting values $Q(0 - 1, 2) = 1$ $Q(\infty - 1, 2) = 0$ (6.4.10) Its relation to the incomplete beta function $I_x(a, b)$ as evaluated by beta1 above is $Q(F - 1, 2) = I_{2/(2+F)}(2, 1/2 + 1/2 F)$ Cumulative Binomial Probability Distribution Suppose an event occurs with probability p per trial. Then the probability P of its occurring k or more times in n trials is termed a cumulative binomial probability, and is related to the incomplete beta function $I_x(a, b)$ as follows: $n P = \sum_{j=k}^n \binom{n}{j} p^j (1-p)^{n-j} = I_p(k, n-k+1)$ (6.4.12) $j = k$

6.5 Bessel Functions of Integer Order 223 For n larger than a dozen or so, beta1 is a much better way to evaluate the sum in (6.4.12) than would be the straightforward sum with concurrent computation of the binomial coefficients. (For n smaller than a dozen, either method is acceptable.) CITED REFERENCES AND FURTHER READING: Abramowitz, M., and Stegun, I.A. 1964, Handbook of Mathematical Functions, Applied Mathematics Series, Volume 55 (Washington: National Bureau of Standards; reprinted 1968 by Dover Publications, New York), Chapters 6 and 26. Pearson, E., and Johnson, N. 1968, Tables of the Incomplete Beta Function (Cambridge:

Cambridge University Press). 6.5 Bessel Functions of Integer Order This section and the next one present practical algorithms for computing various kinds of Bessel functions of integer order. In 6.7 we deal with fractional order. In fact, the more complicated routines for fractional order work fine for integer order too. For integer order, however, the routines in this section (and 6.6) are simpler and faster. Their only drawback is that they are limited by the precision of the underlying rational approximations. For full double precision, it is best to work with the routines for fractional order in 6.7. For any real ν , the Bessel function $J_\nu(x)$ can be defined by the series representation

$$J_\nu(x) = \sum_{k=0}^{\infty} \frac{(-1)^k x^{2k}}{2^k k! \Gamma(\nu + k + 1)} \quad (6.5.1)$$

The series converges for all x , but it is not computationally very useful for $x \gg 1$. For non-integer ν the Bessel function $Y_\nu(x)$ is given by

$$Y_\nu(x) = J_\nu(x) \cos(\nu\pi) - J_{-\nu}(x) \sin(\nu\pi) \quad (6.5.2)$$

The right-hand side goes to the correct limiting value $Y_n(x)$ as ν goes to some integer n , but this is also not computationally useful. For arguments x_j , both Bessel functions look qualitatively like simple power laws, with the asymptotic forms for $0 < x \ll 1$

$$J_\nu(x) \sim \frac{1}{\Gamma(\nu + 1)} x^{\nu} \quad (6.5.3) \quad Y_0(x) \sim -\frac{2}{\pi} \ln(x) \quad Y_\nu(x) \sim -\frac{1}{\Gamma(\nu)} x^{-\nu}$$

6.5 Bessel Functions of Integer Order 223 For n larger than a dozen or so, beta is a much better way to evaluate the sum in (6.4.12) than would be the straightforward sum with concurrent computation of the binomial coefficients. (For n smaller than a dozen, either method is acceptable.) CITED REFERENCES AND FURTHER READING: Abramowitz, M., and Stegun, I.A. 1964, Handbook of Mathematical Functions, Applied Mathematics Series, Volume 55 (Washington: National Bureau of Standards; reprinted 1968 by Dover Publications, New York), Chapters 6 and 26. Pearson, E., and Johnson,

N. 1968, Tables of the Incomplete Beta Function (Cambridge: Cambridge University Press).

6.5 Bessel Functions of Integer Order

This section and the next one present practical algorithms for computing various kinds of Bessel functions of integer order. In 6.7 we deal with fractional order. In fact, the more complicated routines for fractional order work fine for integer order too. For integer order, however, the routines in this section (and 6.6) are simpler and faster. Their only drawback is that they are limited by the precision of the underlying rational approximations. For full double precision, it is best to work with the routines for fractional order in 6.7. For any real ν , the Bessel function $J_\nu(x)$ can be defined by the series representation

$$J_\nu(x) = \sum_{k=0}^{\infty} \frac{(-1)^k x^{\nu+2k}}{2^{\nu+2k} k! \Gamma(\nu+k+1)} \quad (6.5.1)$$

The series converges for all x , but it is not computationally very useful for $x \gg 1$. For non-integer ν the Bessel function $Y_\nu(x)$ is given by

$$Y_\nu(x) = J_\nu(x) \cos(\nu\pi) - J_{-\nu}(x) \sin(\nu\pi) \quad (6.5.2)$$

The right-hand side goes to the correct limiting value $Y_n(x)$ as ν goes to some integer n , but this is also not computationally useful. For arguments x_j , both Bessel functions look qualitatively like simple power laws, with the asymptotic forms for $0 < x_j \ll 1$

$$J_\nu(x) \approx \frac{1}{\Gamma(\nu+1)} \left(\frac{x}{2}\right)^\nu \quad (6.5.3) \quad Y_0(x) \approx -\frac{2}{\pi} \ln(x) \quad Y_\nu(x) \approx -\frac{1}{\Gamma(\nu)} \left(\frac{x}{2}\right)^{-\nu}$$

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Figure 6.5.1. Bessel functions $J_0(x)$ through $J_3(x)$ and $Y_0(x)$ through $Y_2(x)$. For x_j , both Bessel functions look qualitatively like sine or cosine waves whose amplitude decays as $x^{-1/2}$. The asymptotic forms for x are

$$J_\nu(x) \approx \sqrt{\frac{2}{\pi x}} \cos\left(x - \frac{1}{2}\pi\nu - \frac{1}{4}\pi\right) \quad (6.5.4) \quad Y_\nu(x) \approx \sqrt{\frac{2}{\pi x}} \sin\left(x - \frac{1}{2}\pi\nu - \frac{1}{4}\pi\right)$$

In the transition region where $x \sim 1$, the typical amplitudes of the Bessel functions are on the order $1/J(\frac{1}{3}) \approx 2^{1/3} \Gamma(\frac{2}{3}) \approx 0.4473 \sqrt{3} \approx 0.7746$

$$Y_\nu(x) \approx -\frac{2}{\pi} \ln\left(\frac{x}{2}\right) + \frac{1}{\Gamma(\nu)} \left(\frac{x}{2}\right)^{-\nu} \quad (6.5.5)$$

$\sqrt{3} \approx 1.73205$, $\sqrt[3]{-0.7748} \approx -0.9196$, $\sqrt[6]{G(2)} \approx 1.1220$, $\sqrt[3]{3} \approx 1.4422$ which holds asymptotically for large x . Figure 6.5.1 plots the first few Bessel functions of each kind. The Bessel functions satisfy the recurrence relations $J_{n+1}(x) = 2nx J_n(x) - J_{n-1}(x)$ (6.5.6) and $Y_{n+1}(x) = 2nx Y_n(x) - Y_{n-1}(x)$ (6.5.7). As already mentioned in 5.5, only the second of these (6.5.7) is stable in the direction of increasing n for $x \gg n$. The reason that (6.5.6) is unstable in the

direction of increasing n is simply that it is the same recurrence as (6.5.7): A small amount of polluting Y_n introduced by roundoff error will quickly come to swamp the desired J_n , according to equation (6.5.3). A practical strategy for computing the Bessel functions of integer order divides into two tasks: first, how to compute J_0, J_1, Y_0 , and Y_1 , and second, how to use the recurrence relations stably to find other J s and Y s. We treat the first task first: For x between zero and some arbitrary value (we will use the value 8), approximate $J_0(x)$ and $J_1(x)$ by rational functions in x . Likewise approximate by rational functions the regular part of $Y_0(x)$ and $Y_1(x)$, defined as $Y_0(x) = -2/x + J_0(x)\ln(x)$ (6.5.8) and $Y_1(x) = -2/x + J_1(x)\ln(x) - 1$. For $8 \leq x \leq 8$, use the approximating forms $J_n(x) = 2/x \sin(X_n) + P_n/x \cos(X_n) - Q_n$ (6.5.9) and $Y_n(x) = 2/x \cos(X_n) + P_n/x \sin(X_n) + Q_n$ (6.5.10) where $X_n = x - 2n + 1/4$ (6.5.11) and where P_0, P_1, Q_0 , and Q_1 are each polynomials in their arguments, for $0 \leq 8/x \leq 1$. The P s are even polynomials, the Q s odd. Coefficients of the various rational functions and polynomials are given by Hart [1], for various levels of desired accuracy. A straightforward implementation is

```

FUNCTION bessj0(x) REAL bessj0,x
Returns the Bessel function  $J_0(x)$  for any real  $x$ .
REAL ax,xx,z DOUBLE PRECISION p1,p2,p3,p4,p5,q1,q2,q3,q4,q5,r1,r2,r3,r4, *
r5,r6,s1,s2,s3,s4,s5,s6,y
Well accumulate polynomials in double

```

```

precision. SAVE p1,p2,p3,p4,p5,q1,q2,q3,q4,q5,r1,r2,r3,r4,r5,r6,
* s1,s2,s3,s4,s5,s6 DATA p1,p2,p3,p4,p5/1.d0,-.1098628627d-
2,.2734510407d-4, * -.2073370639d-5,.2093887211d-6/, q1,q2,q3,q4,q5/-
.1562499995d-1, * .1430488765d-3,-.6911147651d-5,.7621095161d-
6,-.934945152d-7/ DATA r1,r2,r3,r4,r5,r6/57568490574.d0,-13362590354.
* -11214424.18d0,77392.33017d0,-184.9052456d0/, * s1,s2,s3,s4,s5,s6/575
* 9494680.718d0,59272.64853d0,267.8532712d0,1.d0/ if(abs(x).lt.8.)then
Direct rational function fit.  $y=x^{**2}$  bessj0=(r1+y*(r2+y*(r3+y*(r4+y*
* /(s1+y*(s2+y*(s3+y*(s4+y*(s5+y*s6)))))) else Fitting func-
tion (6.5.9).  $ax=abs(x)$   $z=8./ax$   $y=z^{**2}$   $xx=ax-.785398164$ 
bessj0=sqrt(.636619772/ax)*(cos(xx)*(p1+y*(p2+y*(p3+y*(p4+y
226 Chapter 6. Special Functions * *p5))))-z*sin(xx)*(q1+y*(q2+y*
endif return END FUNCTION bessy0(x) REAL bessy0,x C
USES bessj0 Returns the Bessel function Y0 (x) for positive x.
REAL xx,z,bessj0 DOUBLE PRECISION p1,p2,p3,p4,p5,q1, *
q2,q3,q4,q5,r1,r2,r3,r4, * r5,r6,s1,s2,s3,s4,s5,s6,y Wel l accumu-
late polynomials in double precision. SAVE p1,p2,p3,p4,p5,q1,q2,q3,q4,q5,
* r5,r6,s1,s2,s3,s4,s5,s6 DATA p1,p2,p3,p4,p5/1.d0,-.1098628627d-
2,.2734510407d-4, * -.2073370639d-5,.2093887211d-6/, q1,q2,q3,q4,q5/-
.1562499995d-1, * .1430488765d-3,-.6911147651d-5,.7621095161d-
6,-.934945152d-7/ DATA r1,r2,r3,r4,r5,r6/-2957821389.d0,7062834065.d0
512359803.6d0, * 10879881.29d0,-86327.92757d0,228.4622733d0/,
* s1,s2,s3,s4,s5,s6/40076544269.d0,745249964.8d0, * 7189466.438d0,4744
if(x.lt.8.)then Rational function approximation of (6.5.8).  $y=x^{**2}$ 
bessy0=(r1+y*(r2+y*(r3+y*(r4+y*(r5+y*r6))))/(s1+y*(s2+y
* *(s3+y*(s4+y*(s5+y*s6))))+.636619772*bessj0(x)*log(x) else
Fitting function (6.5.10).  $z=8./x$   $y=z^{**2}$   $xx=x-.785398164$ 
bessy0=sqrt(.636619772/x)*(sin(xx)*(p1+y*(p2+y*(p3+y*(p4+y*
* p5))))+z*cos(xx)*(q1+y*(q2+y*(q3+y*(q4+y*q5)))) endif
return END FUNCTION bessj1(x) REAL bessj1,x Returns the
Bessel function J1 (x) for any real x. REAL ax,xx,z DOU-
BLE PRECISION p1,p2,p3,p4,p5,q1,q2,q3,q4,q5,r1,r2,r3,r4, *
```

```

r5,r6,s1,s2,s3,s4,s5,s6,y Well accumulate polynomials in double
precision. SAVE p1,p2,p3,p4,p5,q1,q2,q3,q4,q5,r1,r2,r3,r4,r5,r6,
* s1,s2,s3,s4,s5,s6 DATA r1,r2,r3,r4,r5,r6/72362614232.d0,-7895059235.d0
* -2972611.439d0,15704.48260d0,-30.16036606d0/, * s1,s2,s3,s4,s5,s6/144
* 18583304.74d0,99447.43394d0,376.9991397d0,1.d0/ DATA p1,p2,p3,p4,
2,-.3516396496d-4,.2457520174d-5, * -.240337019d-6/, q1,q2,q3,q4,q5/.04
.2002690873d-3, * .8449199096d-5,-.88228987d-6,.105787412d-
6/ if(abs(x).lt.8.)then Direct rational approximation. y=x**2
bessj1=x*(r1+y*(r2+y*(r3+y*(r4+y*(r5+y*r6)))) * /(s1+y*(s2+y*(s3
else Fitting function (6.5.9). ax=abs(x) z=8./ax y=z**2
xx=ax-2.356194491

```

6.5 Bessel Functions of Integer Order 227 $\text{bessj1} = \sqrt{.636619772/\text{ax}}$

```

* *p5))))-z*sin(xx)*(q1+y*(q2+y*(q3+y*(q4+y*q5)))) * *sign(1.,x)
endif return END FUNCTION bessy1(x) REAL bessy1,x C
USES bessj1 Returns the Bessel function Y1 (x) for positive x.
REAL xx,z,bessj1 DOUBLE PRECISION p1,p2,p3,p4,p5,q1,q2,q3,q4,q5,
* r5,r6,s1,s2,s3,s4,s5,s6,s7,y Well accumulate polynomials in
double precision. SAVE p1,p2,p3,p4,p5,q1,q2,q3,q4,q5,r1,r2,r3,r4,
* r5,r6,s1,s2,s3,s4,s5,s6,s7 DATA p1,p2,p3,p4,p5/1.d0,.183105d-
2,-.3516396496d-4,.2457520174d-5, * -.240337019d-6/, q1,q2,q3,q4,q5/.04
.2002690873d-3, * .8449199096d-5,-.88228987d-6,.105787412d-
6/ DATA r1,r2,r3,r4,r5,r6/-.4900604943d13,.1275274390d13,-
.5153438139d11, * .7349264551d9,-.4237922726d7,.8511937935d4/,
* s1,s2,s3,s4,s5,s6,s7/.2499580570d14,.4244419664d12, * .3733650367d10,
if(x.lt.8.)then Rational function approximation of (6.5.8). y=x**2
bessy1=x*(r1+y*(r2+y*(r3+y*(r4+y*(r5+y*r6))))/(s1+y*(s2+y*
* (s3+y*(s4+y*(s5+y*(s6+y*s7))))))+.636619772 * *(bessj1(x)*log(x)-
1./x) else Fitting function (6.5.10). z=8./x y=z**2 xx=x-
2.356194491 bessy1=sqrt(.636619772/x)*(sin(xx)*(p1+y*(p2+y*(p3+y*
* *p5))))+z*cos(xx)*(q1+y*(q2+y*(q3+y*(q4+y*q5)))) en-
dif return END We now turn to the second task, namely how to
use the recurrence formulas (6.5.6) and (6.5.7) to get the Bessel

```

functions $J_n(x)$ and $Y_n(x)$ for $n = 2$. The latter of these is straightforward, since its upward recurrence is always stable:

```

FUNCTION bessy(n,x) INTEGER n REAL bessy,x C USES
bessy0,bessy1 Returns the Bessel function  $Y_n(x)$  for positive x
and  $n = 2$ . INTEGER j REAL by,bym,byp,tox,bessy0,bessy1
if(n.lt.2)pause bad argument n in bessy tox=2./x by=bessy1(x)
Starting values for the recurrence. bym=bessy0(x) do 1
1 j=1,n-1 Recurrence (6.5.7). byp=j*tox*by-bym bym=by
by=byp enddo 1 1 bessy=by return END

```

6.6 Modified Bessel Functions of Integer Order 229 $b_j = b_j * \text{BIGNI}$
 $b_{j+1} = b_{j+1} * \text{BIGNI}$ $sum = sum * \text{BIGNI}$ endif if(jsum.ne.0) $sum = sum + b_j$
Accumulate the sum. $jsum = 1 - jsum$ Change 0 to 1 or vice versa.
if(j.eq.n) $b_{j+1} = b_j$ Save the unnormalized answer. enddo 1 2
 $sum = 2. * sum - b_j$ Compute (5.5.16) $b_{j+1} = b_{j+1} / sum$ and use it
to normalize the answer. endif if(x.lt.0..and.mod(n,2).eq.1) $b_{j+1} = -$
 b_{j+1} return END CITED REFERENCES AND FURTHER

READING: Abramowitz, M., and Stegun, I.A. 1964, Handbook of Mathematical Functions, Applied Mathematics Series, Volume 55 (Washington: National Bureau of Standards; reprinted 1968 by Dover Publications, New York), Chapter 9.
Hart, J.F., et al. 1968, Computer Approximations (New York: Wiley), 6.8, p. 141. [1] 6.6 Modified Bessel Functions of Integer Order The modified Bessel functions $I_n(x)$ and $K_n(x)$ are equivalent to the usual Bessel functions J_n and Y_n evaluated for purely imaginary arguments. In detail, the relationship is
 $I_n(x) = (-i)^n J_n(ix)$ (6.6.1) $K_n(x) = \frac{1}{2} \int_0^\infty [J_n(ix) + iY_n(ix)]$
The particular choice of prefactor and of the linear combination of J_n and Y_n to form K_n are simply choices that make the functions real-valued for real arguments x . For small arguments $x \ll n$, both $I_n(x)$ and $K_n(x)$ become, asymptotically, simple powers of their argument
 $n = 0$ $x I_n(x) \approx \frac{1}{n!} x^n$ (6.6.2)
 $K_0(x) \approx -\ln(x)$ $-n \ll 0$ $x K_n(x) \approx \frac{(n-1)!}{2} x^{-n}$ These expressions

are virtually identical to those for $J_n(x)$ and $Y_n(x)$ in this region, except for the factor of $-2/p$ difference between $Y_n(x)$ and $K_n(x)$. In the region

230 Chapter 6. Special Functions I0 K0 K1 K2 I1 I2 I3
0 1 2 3 4 01234 x Figure 6.6.1. Modified Bessel functions I0(x) through I3(x), K0(x) through K2(x). x n, however, the modified functions have quite different behavior than the Bessel functions, $v I_n(x) \approx \frac{1}{2} p x \exp(x)$ (6.6.3) $v 2 p x \exp(-x) K_n(x) \approx p$ The modified functions evidently have exponential rather than sinusoidal behavior for large arguments (see Figure 6.6.1). The smoothness of the modified Bessel functions, once the exponential factor is removed, makes a simple polynomial approximation of a few terms quite suitable for the functions I0, I1, K0, and K1. The following routines, based on polynomial coefficients given by Abramowitz and Stegun [1], evaluate these four functions, and will provide the basis for upward recursion for $n \geq 1$ when $x \geq n$.
FUNCTION bessio(x) REAL bessio,x Returns the modified Bessel function I0(x) for any real x. REAL ax DOUBLE PRECISION p1,p2,p3,p4,p5,p6,p7,q1,q2,q3,q4,q5,q6,q7,* q8,q9,y Accumulate polynomials in double precision. SAVE p1,p2,p3,p4,p5,p6,p7,q1,q2,q3,q4,q5,q6,q7,q8,q9 DATA p1,p2,p3,p4,p5,p6,p7,q1,q2,q3,q4,q5,q6,q7,q8,q9
* 0.2659732d0,0.360768d-1,0.45813d-2/ DATA q1,q2,q3,q4,q5,q6,q7,q8,q9
1, * 0.225319d-2,-0.157565d-2,0.916281d-2,-0.2057706d-1, * 0.2635537d-
1,-0.1647633d-1,0.392377d-2/

6.6 Modified Bessel Functions of Integer Order 231 if
(abs(x).lt.3.75) then y=(x/3.75)**2 bessio=p1+y*(p2+y*(p3+y*(p4+y*
else ax=abs(x) y=3.75/ax bessio=(exp(ax)/sqrt(ax))*(q1+y*(q2+y*(q3*
* +y*(q5+y*(q6+y*(q7+y*(q8+y*q9)))))) endif return END

FUNCTION bessk0(x) REAL bessk0,x C USES bessio Re-
turns the modified Bessel function K0(x) for positive real x.
REAL bessio DOUBLE PRECISION p1,p2,p3,p4,p5,p6,p7,q1,

```

* q2,q3,q4,q5,q6,q7,y Accumulate polynomials in double precision.
SAVE p1,p2,p3,p4,p5,p6,p7,q1,q2,q3,q4,q5,q6,q7 DATA
p1,p2,p3,p4,p5,p6,p7/-0.57721566d0,0.42278420d0,0.23069756d0,
* 0.3488590d-1,0.262698d-2,0.10750d-3,0.74d-5/ DATA q1,q2,q3,q4,q5,q6
0.7832358d-1,0.2189568d-1, * -0.1062446d-1,0.587872d-2,-0.251540d-
2,0.53208d-3/ if (x.le.2.0) then Polynomial fit. y=x*x/4.0
bessk0=(-log(x/2.0)*bessi0(x))+(p1+y*(p2+y*(p3+ * y*(p4+y*(p5+y*
else y=(2.0/x) bessk0=(exp(-x)/sqrt(x))*(q1+y*(q2+y*(q3+
* y*(q4+y*(q5+y*(q6+y*q7)))))) endif return END FUNC-
TION bessi1(x) REAL bessi1,x Returns the modified Bessel
function I1 (x) for any real x. REAL ax DOUBLE PRECI-
SION p1,p2,p3,p4,p5,p6,p7,q1,q2,q3,q4,q5,q6,q7, * q8,q9,y Ac-
cumulate polynomials in double precision. SAVE p1,p2,p3,p4,p5,p6,p7,q1
DATA p1,p2,p3,p4,p5,p6,p7/0.5d0,0.87890594d0,0.51498869d0,
* 0.15084934d0,0.2658733d-1,0.301532d-2,0.32411d-3/ DATA
q1,q2,q3,q4,q5,q6,q7,q8,q9/0.39894228d0,-0.3988024d-1, * -0.362018d-
2,0.163801d-2,-0.1031555d-1,0.2282967d-1, * -0.2895312d-1,0.1787654d-
1,-0.420059d-2/ if (abs(x).lt.3.75) then Polynomial fit. y=(x/3.75)**2
bessi1=x*(p1+y*(p2+y*(p3+y*(p4+y*(p5+y*(p6+y*p7))))))
else ax=abs(x) y=3.75/ax bessi1=(exp(ax)/sqrt(ax))*(q1+y*(q2+y*(q3
* y*(q5+y*(q6+y*(q7+y*(q8+y*q9)))))) if(x.lt.0.)bessi1=-
bessi1 endif return END

```

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```

FUNCTION bessk1(x)
REAL bessk1,x C USES bessi1 Returns the modified Bessel
function K1 (x) for positive real x. REAL bessi1 DOUBLE
PRECISION p1,p2,p3,p4,p5,p6,p7,q1, * q2,q3,q4,q5,q6,q7,y
Accumulate polynomials in double precision. SAVE p1,p2,p3,p4,p5,p6,p7
DATA p1,p2,p3,p4,p5,p6,p7/1.0d0,0.15443144d0,-0.67278579d0,
* -0.18156897d0,-0.1919402d-1,-0.110404d-2,-0.4686d-4/ DATA
q1,q2,q3,q4,q5,q6,q7/1.25331414d0,0.23498619d0,-0.3655620d-1,
* 0.1504268d-1,-0.780353d-2,0.325614d-2,-0.68245d-3/ if (x.le.2.0)
then Polynomial fit. y=x*x/4.0 bessk1=(log(x/2.0)*bessi1(x))+(1.0/x)*

```

```

* y*(p3+y*(p4+y*(p5+y*(p6+y*p7)))))) else y=2.0/x bessk1=(exp(-
x)/sqrt(x))*(q1+y*(q2+y*(q3+ * y*(q4+y*(q5+y*(q6+y*q7))))))
endif return END

```

The recurrence relation for $I_n(x)$ and $K_n(x)$ is the same as that for $J_n(x)$ and $Y_n(x)$ provided that ix is substituted for x . This has the effect of changing a sign in the relation, $I_{n+1}(x) = -2n/x I_n(x) + I_{n-1}(x)$ (6.6.4) $K_{n+1}(x) = +2n/x K_n(x) + K_{n-1}(x)$. These relations are always unstable for upward recurrence. For K_n , itself growing, this presents no problem. For I_n , however, the strategy of downward recursion is therefore required once again, and the starting point for the recursion may be chosen in the same manner as for the routine `bessj`. The only fundamental difference is that the normalization formula for $I_n(x)$ has an alternating minus sign in successive terms, which again arises from the substitution of ix for x in the formula used previously for $J_{n-1} = I_0(x) - 2I_2(x) + 2I_4(x) - 2I_6(x) + \dots$ (6.6.5). In fact, we prefer simply to normalize with a call to `bessi0`. With this simple modification, the recursion routines `bessj` and `bessy` become the new routines `bessi` and `bessk`:

```

FUNCTION bessk(n,x) INTEGER
n REAL bessk,x C USES bessk0,bessk1 Returns the modified
Bessel function  $K_n(x)$  for positive  $x$  and  $n = 2$ . INTEGER j
REAL bk,bkm,bkp,tox,bessk0,bessk1 if (n.lt.2) pause bad ar-
gument n in bessk tox=2.0/x

```

6.6 Modified Bessel Functions of Integer Order 233 `bkm=bessk0(x)`

```

Upward recurrence for all x... bk=bessk1(x) do 1 1 j=1,n-
1 ...and here it is. bkp=bkm+j*tox*bk bkm=bk bk=bkp
enddo 1 1 bessk=bk return END
FUNCTION bessi(n,x) IN-
TEGER n,IACC REAL bessi,x,BIGNO,BIGNI PARAMETER
(IACC=40,BIGNO=1.0e10,BIGNI=1.0e-10) C USES bessi0
Returns the modified Bessel function  $I_n(x)$  for any real  $x$  and
 $n = 2$ . INTEGER j,m REAL bi,bim,bip,tox,bessi0 if (n.lt.2)
pause bad argument n in bessi if (x.eq.0.) then bessi=0. else

```

```

tox=2.0/abs(x) bip=0.0 bi=1.0 beessi=0. m=2*((n+int(sqrt(float(IACC*
Downward recurrence from even m. do 1 1 j=m,1,-1 Make
IACC larger to increase accuracy. bim=bip+float(j)*tox*bi
The downward recurrence. bip=bi bi=bim if (abs(bi).gt.BIGNO)
then Renormalize to prevent over ows. beessi=beessi*BIGNI
bi=bi*BIGNI bip=bip*BIGNI endif if (j.eq.n) beessi=bip enddo
1 1 beessi=beessi*beessi0(x)/bi Normalize with beessi0. if (x.lt.0..and.mod(n,
beessi=-beessi endif return END CITED REFERENCES AND
FURTHER READING: Abramowitz, M., and Stegun, I.A.
1964, Handbook of Mathematical Functions, Applied Mathe-
matics Series, Volume 55 (Washington: National Bureau of
Standards; reprinted 1968 by Dover Publications, New York),
9.8. [1] Carrier, G.F., Krook, M. and Pearson, C.E. 1966, Func-
tions of a Complex Variable (New York: McGraw-Hill), pp.
220ff.

```

234 Chapter 6. Special Functions 6.7 Bessel Functions of Fractional Order, Airy Functions, Spherical Bessel Functions

Many algorithms have been proposed for computing Bessel functions of fractional order numerically. Most of them are, in fact, not very good in practice. The routines given here are rather complicated, but they can be recommended wholeheartedly.

Ordinary Bessel Functions The basic idea is Steeds method, which was originally developed [1] for Coulomb wave functions. The method calculates J_ν , $J_{-\nu}$, Y_ν , and $Y_{-\nu}$ simultaneously, and so involves four relations among these functions. Three of the relations come from two continued fractions, one of which is complex. The fourth is provided by the Wronskian relation $W = J_\nu Y_{-\nu} - Y_\nu J_{-\nu} = 2/\pi x$ (6.7.1)

The first continued fraction, CF1, is defined by $f = J_\nu / J_{\nu+1} = x - J_{\nu+1} / (J_\nu + 1)$ (6.7.2)

$$1 = x - \frac{1}{2} \frac{(\nu+1)}{x} - \frac{2}{x} \frac{(\nu+2)}{x} - \dots$$

You can easily derive it from the three-term recurrence relation for Bessel functions: Start with equation (6.5.6) and use equation (5.5.18). Forward eval-

uation of the continued fraction by one of the methods of 5.2 is essentially equivalent to backward recurrence of the recurrence relation. The rate of convergence of CF1 is determined by the position of the turning point $x_t = p + 1$, beyond which the Bessel functions become oscillatory. If $x \ll x_t$, convergence is very rapid. If $x \gg x_t$, then each iteration of the continued fraction effectively increases by one until $x \ll x_t$; thereafter rapid convergence sets in. Thus the number of iterations of CF1 is of order x for large x . In the routine `bessjy` we set the maximum allowed number of iterations to 10,000. For larger x , you can use the usual asymptotic expressions for Bessel functions. One can show that the sign of J is the same as the sign of the denominator of CF1 once it has converged. The complex continued fraction CF2 is defined by $(1/2)^2 - 2(3/2)^2 - 2 = -1$ $p + iq = J + iY$ $2x + i + i x^2(x + i) + 2(x + 2i) +$ (6.7.3) $J + iY$ (We sketch the derivation of CF2 in the analogous case of modified Bessel functions in the next subsection.) This continued fraction converges rapidly for $x \gg x_t$, while convergence fails as $x \rightarrow 0$. We have to adopt a special method for small x , which we describe below. For x not too small, we can ensure that $x \gg x_t$ by a stable recurrence of J and J' downwards to a value $\approx 1/x$, thus yielding the ratio f at this lower value of x . This is the stable direction for the recurrence relation. The initial values for the recurrence are $J = \text{arbitrary}, J' = f J$, (6.7.4) with the sign of the arbitrary initial value of J chosen to be the sign of the denominator of CF1. Choosing the initial value of J very small minimizes the possibility of overflow during the recurrence. The recurrence relations are $J_{-1} = xJ + J'$ (6.7.5) $J_{-1} = -1/x J_{-1} - J$

6.7 Bessel Functions of Fractional Order 235 Once CF2 has been evaluated at x , then with the Wronskian (6.7.1) we have enough relations to solve for all four quantities. The formulas

are simplified by introducing the quantity $(6.7.6) = p - f q$ Then $1/2$ (6.7.7) $J = W q + (p - f)$ (6.7.8) $J = f J$ (6.7.9) $Y = J$ (6.7.10) $Y = Y p + q$

The sign of J in (6.7.7) is chosen to be the same as the sign of the initial J in (6.7.4). Once all four functions have been determined at the value $=$, we can find them at the original value of $.$ For J and J , simply scale the values in (6.7.4) by the ratio of (6.7.7) to the value found after applying the recurrence (6.7.5). The quantities Y and Y can be found by starting with the values in (6.7.9) and (6.7.10) and using the stable upwards recurrence $Y_{+1} = 2 x Y - Y_{-1}$ (6.7.11) together with the relation (6.7.12) $Y = xY - Y_{+1}$ Now turn to the case of small x , when CF2 is not suitable. Temme [2] has given a good method of evaluating Y and Y_{+1} , and hence Y from (6.7.12), by series expansions that accurately handle the singularity as $x \rightarrow 0$. The expansions work only for $— = 1/2$, and so now the recurrence (6.7.5) is used to evaluate f at a value $=$ in this interval. Then one calculates J from (6.7.13) $J = W Y - Y$ f and J from (6.7.8). The values at the original value of are determined by scaling as before, and the Y s are recurred up as before. Temmes series are $8 8 Y = - c_k g_k Y_{+1} = - 2 c_k h_k$ (6.7.14) $x^k = 0 k = 0$ Here (6.7.15) $c_k = (-x^2 / 4)^k k!$ while the coefficients g_k and h_k are defined in terms of quantities p_k , q_k , and f_k that can be found by recursion: $g_k = f_k + 2 \sin^2 p$ $2 q_k h_k = -k g_k + p_k$ (6.7.16) $p_k = p_{k-1} k - q_k = q_{k-1} k + f_k = k f_{k-1} + p_{k-1} + q_{k-1} k^2 - 2$

236 Chapter 6. Special Functions The initial values for the recurrences are $- G(1 +) x p_0 = 1 p^2 G(1 -) x$ (6.7.17) $q_0 = 1 p^2 p f_0 = 2 p s \ln 2 \sin p \cosh s G_1 () + \sinh s x G_2 ()$ with $s = \ln 2 x 1$ (6.7.18) $G_1 () = 1 2 G(1 -) - 1 G(1 +) 1 G_2 () = 1 2 G(1 -) + 1 G(1 +)$ The whole point of writing the formulas in this way is that the potential problems as 0

can be controlled by evaluating $p/\sin p$, $\sinh s/s$, and $G1$ carefully. In particular, Temme gives Chebyshev expansions for $G1$ () and $G2$ (). We have rearranged his expansion for $G1$ to be explicitly an even series in x so that we can use our routine `chebev` as explained in 5.8. The routine assumes $\nu = 0$. For negative ν you can use the reflection formulas $J_{-\nu} = \cos \nu J_{\nu} - \sin \nu Y_{\nu}$ (6.7.19) $Y_{-\nu} = \sin \nu J_{\nu} + \cos \nu Y_{\nu}$. The routine also assumes $x \neq 0$. For $x \neq 0$ the functions are in general complex, but expressible in terms of functions with $x \neq 0$. For $x = 0$, Y_{ν} is singular. Internal arithmetic in the routine is carried out in double precision. To maintain portability, complex arithmetic has been recoded with real variables.

SUBROUTINE `bessjy(x,xnu,rj,ry,rjp,ryp)`
 INTEGER MAXIT REAL rj,rjp,ry,ryp,x,xnu,XMIN DOUBLE PRECISION EPS,FPMIN,PI
 PARAMETER (EPS=1.e-10,FPMIN=1.e-30,MAXIT=10000,XMIN=2., * PI=3.141592653589793d0)

C USES `beschb` Returns the Bessel functions $rj = J_{\nu}$, $ry = Y_{\nu}$ and their derivatives $rjp = J'_{\nu}$, $ryp = Y'_{\nu}$, for positive x and for $xnu = \nu = 0$. The relative accuracy is within one or two significant digits of EPS, except near a zero of one of the functions, where EPS controls its absolute accuracy. FPMIN is a number close to the machines smallest floating-point number. All internal arithmetic is in double precision. To convert the entire routine to double precision, change the REAL declaration above and decrease EPS to 10^{-16} . Also convert the subroutine `beschb`.

INTEGER i,sign,l,nl DOUBLE PRECISION a,b,br,bi,c,cr,ci,d,del,del1,den,di,dlr,dli, * dr,e,f,fact,fact2,fact3,ff,gam,gam1,gam2,gammi,gampl,h, * p,pimu,pimu2 * rymu,rymup,rytemp,sum,sum1,temp,w,x2,xi,xi2,xmu,xmu2
 if(x.le.0..or pause bad arguments in `bessjy` if(x.lt.XMIN) then nl is the number of downward recurrences of the J s and upward recurrences of the Y s. xmu lies between $-1/2$ and $1/2$ for $x \geq XMIN$, while it is chosen so that x is greater than the turning point for x

= XMIN. nl=int(xnu+.5d0) else nl=max(0,int(xnu-x+1.5d0))
 endif xmu=xnu-nl xmu2=xmu*xmu xi=1.d0/x xi2=2.d0*xi
 w=xi2/PI The Wronskian.

6.7 Bessel Functions of Fractional Order 237 isign=1 Evaluate CF1 by modified Lentzs method (5.2). isign keeps track of sign changes in the denominator. h=xnu*xi if(h.lt.FPMIN)h=FPMIN
 b=xi2*xnu d=0.d0 c=h do 1 1 i=1,MAXIT b=b+xi2 d=b-d
 if(abs(d).lt.FPMIN)d=FPMIN c=b-1.d0/c if(abs(c).lt.FPMIN)c=FPMIN
 d=1.d0/d del=c*d h=del*h if(d.lt.0.d0)isign=-isign if(abs(del-1.d0).lt.EPS)goto 1 enddo 1 1 pause x too large in bessjy;
 try asymptotic expansion 1 continue rjl=isign*FPMIN Initialize J and J for downward recurrence. rjpl=h*rjl rjl1=rjl
 Store values for later rescaling. rjp1=rjpl fact=xnu*xi do 1 2
 l=nl,1,-1 rjtemp=fact*rjl+rjpl fact=fact-xi rjpl=fact*rjtemp
 rjl rjl=rjtemp enddo 1 2 if(rjl.eq.0.d0)rjl=EPS f=rjpl/rjl Now
 have unnormalized J and J . if(x.lt.XMIN) then Use series.
 x2=.5d0*x pimu=PI*xmu if(abs(pimu).lt.EPS)then fact=1.d0
 else fact=pimu/sin(pimu) endif d=-log(x2) e=xmu*d if(abs(e).lt.EPS)then
 fact2=1.d0 else fact2=sinh(e)/e endif call beschb(xmu,gam1,gam2,gampl
 Chebyshev evaluation of G1 and G2 . ff=2.d0/PI*fact*(gam1*cosh(e)+ga
 f0 . e=exp(e) p=e/(gampl*PI) p0 . q=1.d0/(e*PI*gammi)
 q0 . pimu2=0.5d0*pimu if(abs(pimu2).lt.EPS)then fact3=1.d0
 else fact3=sin(pimu2)/pimu2 endif r=PI*pimu2*fact3*fact3
 c=1.d0 d=-x2*x2 sum=ff+r*q sum1=p

238 Chapter 6. Special Functions do 1 3 i=1,MAXIT
 ff=(i*ff+p+q)/(i*i-xmu2) c=c*d/i p=p/(i-xmu) q=q/(i+xmu)
 del=c*(ff+r*q) sum=sum+del del1=c*p-i*del sum1=sum1+del1
 if(abs(del).lt.(1.d0+abs(sum))*EPS)goto 2 enddo 1 3 pause
 bessy series failed to converge 2 continue rymu=-sum ry1=-
 sum1*xi2 rymup=xmu*xi*rymu-ry1 rjmu=w/(rymup-f*rymu)
 Equation (6.7.13). else Evaluate CF2 by modified Lentzs
 method (5.2). a=.25d0-xmu2 p=-.5d0*xi q=1.d0 br=2.d0*x


```

bi=2.d0 fact=a*xi/(p*p+q*q) cr=br+q*fact ci=bi+p*fact den=br*br+br*dr+dr*dr
dr=br/den di=-bi/den dlr=cr*dr-ci*di dli=cr*di+ci*dr temp=p*dli-q*dli
q=dli+q*dli p=temp do 1 4 i=2,MAXIT a=a+2*(i-1)
bi=bi+2.d0 dr=a*dr+br di=a*di+bi if(abs(dr)+abs(di).lt.FPMIN)dr=F
fact=a/(cr*cr+ci*ci) cr=br+cr*fact ci=bi-ci*fact if(abs(cr)+abs(ci).lt.F
den=dr*dr+di*di dr=dr/den di=-di/den dlr=cr*dr-ci*di dli=cr*di+ci*d
temp=p*dli-q*dli q=p*dli+q*dli p=temp if(abs(dli-1.d0)+abs(dli).lt.EF
3 enddo 1 4 pause cf2 failed in bessjy 3 continue gam=(p-f)/q
Equations (6.7.6) (6.7.10). rjmu=sqrt(w/((p-f)*gam+q))
rjmu=sign(rjmu,rjl) rymu=rjmu*gam rymup=rymu*(p+q/gam)
ryl=xmu*xi*rymu-rymup endif fact=rjmu/rjl

```

246 Chapter 6. Special Functions 6.8 Spherical Harmonics

Spherical harmonics occur in a large variety of physical problems, for example, whenever a wave equation, or Laplace's equation, is solved by separation of variables in spherical coordinates. The spherical harmonic $Y_l^m(\theta, \phi)$, $-l \leq m \leq l$, is a function of the two coordinates θ, ϕ on the surface of a sphere. The spherical harmonics are orthogonal for different l and m , and they are normalized so that their integrated square over the sphere is unity:

$$\int_0^{2\pi} \int_0^\pi \sin \theta Y_l^m(\theta, \phi) Y_{l'}^{m'}(\theta, \phi) d\theta d\phi = \delta_{ll'} \delta_{mm'} \quad (6.8.1)$$

Here asterisk denotes complex conjugation. Mathematically, the spherical harmonics are related to associated Legendre polynomials by the equation

$$Y_l^m(\theta, \phi) = \sqrt{\frac{2l+1}{4\pi} \frac{(l-m)!}{(l+m)!}} (\cos \theta)^m P_l^m(\cos \theta) e^{im\phi} \quad (6.8.2)$$

By using the relation $Y_l^{-m}(\theta, \phi) = (-1)^m Y_l^m(\theta, \phi)^*$ (6.8.3) we can always relate a spherical harmonic to an associated Legendre polynomial with $m \geq 0$. With $x = \cos \theta$, these are defined in terms of the ordinary Legendre polynomials (cf. 4.5 and 5.5) by

$$P_l^m(x) = (-1)^m (1-x^2)^{m/2} \frac{d^m}{dx^m} P_l(x) \quad (6.8.4)$$

The first few associated Legendre polynomials, and their corresponding normalized spherical harmonics, are

$$P_0^0(x) = 1 \quad Y_0^0 = \frac{1}{\sqrt{4\pi}}$$

$$P_1^0(x) = x \quad Y_1^0 = \sqrt{\frac{3}{4\pi}} \cos \theta$$

$$P_1^1(x) = \sqrt{1-x^2} \quad Y_1^1 = \sqrt{\frac{3}{8\pi}} \sin \theta e^{i\phi}$$

$$Y_{l,m}^2 = \frac{(-1)^m (1+m)!}{2^m m! (1-m)!} (1-x^2)^{m/2} P_{l-m}^{(m,m)}(x)$$

$$Y_{l,m}^1 = -\frac{1}{2} x Y_{l,m}^2 - \frac{(1-m)(m+1)}{2(l-m)(l+m)} Y_{l,m-1}^2$$
 (6.8.5) There are many bad ways to evaluate associated Legendre polynomials numerically. For example, there are explicit expressions, such as

247 6.8 Spherical Harmonics where the polynomial continues up through the term in $(1-x)^{l-m}$. (See [1] for this and related formulas.) This is not a satisfactory method because evaluation of the polynomial involves delicate cancellations between successive terms, which alternate in sign. For large l , the individual terms in the polynomial become very much larger than their sum, and all accuracy is lost. In practice, (6.8.6) can be used only in single precision (32-bit) for l up to 6 or 8, and in double precision (64-bit) for l up to 15 or 18, depending on the precision required for the answer. A more robust computational procedure is therefore desirable, as follows: The associated Legendre functions satisfy numerous recurrence relations, tabulated in [1-2]. These are recurrences on l alone, on m alone, and on both l and m simultaneously. Most of the recurrences involving m are unstable, and so dangerous for numerical work. The following recurrence on l is, however, stable (compare 5.5.1):

$$(l-m)P_{l,m}^1 - x(2l-1)P_{l,m-1}^1 - (l+m-1)P_{l,m}^1 = 0$$

(6.8.7) It is useful because there is a closed-form expression for the starting value,

$$P_{l,m}^m = (-1)^m (2m-1)!! (1-x^2)^{m/2}$$

(6.8.8) (The notation $n!!$ denotes the product of all odd integers less than or equal to n .) Using (6.8.7) with $l = m+1$, and setting $P_{m,m}^{m-1} = 0$, we find

$$P_{m,m}^m = x(2m+1)P_{m,m}^{m-1}$$

(6.8.9) Equations (6.8.8) and (6.8.9) provide the two starting values required for (6.8.7)

for general l . The function that implements this is FUNCTION `plgndr(l,m,x)` INTEGER `l,m` REAL `plgndr,x` Computes the associated Legendre polynomial $P_m^l(x)$. Here `m` and `l` are integers satisfying $0 \leq m \leq l$, while `x` lies in the range $-1 \leq x \leq 1$. INTEGER `i,ll` REAL `fact,pll,pmm,pmmp1,somx2` if(`m.lt.0.or.m.gt.l.or.abs(x).gt.1.`)pause bad arguments in `plgndr` `pmm=1`. Compute P_m^m . if(`m.gt.0`) then `somx2=sqrt((1-x)*(1+x))` `fact=1`. do 1 1 `i=1,m` `pmm=-pmm*fact*somx2` `fact=fact+2`. enddo 1 1 endif if(`l.eq.m`) then `plgndr=pmm` else `pmmp1=x*(2*m+1)*pmm` Compute P_m^{m+1} . if(`l.eq.m+1`) then `plgndr=pmmp1` else Compute $P_m^{l,m+1}$. do 1 2 `ll=m+2,l`

248 Chapter 6. Special Functions `pll=(x*(2*ll-1)*pmmp1-(ll+m-1)*pmm)/(ll-m)` `pmm=pmmp1` `pmmp1=pll` enddo 1 2 `plgndr=pll` endif endif return END CITED REFERENCES AND FURTHER READING: Magnus, W., and Oberhettinger, F. 1949, *Formulas and Theorems for the Functions of Mathematical Physics* (New York: Chelsea), pp. 54ff. [1] Abramowitz, M., and Stegun, I.A. 1964, *Handbook of Mathematical Functions*, Applied Mathematics Series, Volume 55 (Washington: National Bureau of Standards; reprinted 1968 by Dover Publications, New York), Chapter 8. [2] 6.9 Fresnel Integrals, Cosine and Sine Integrals Fresnel Integrals The two Fresnel integrals are defined by $\cos p \sin p$ $C(x) = \int_0^x \cos t^2 dt$ (6.9.1) $S(x) = \int_0^x \sin t^2 dt$ The most convenient way of evaluating these functions to arbitrary precision is to use power series for small x and a continued fraction for large x . The series are $2 x^5 4 x^9$ $C(x) = x - \frac{p^2}{5} \frac{2!}{2!} + \frac{p^2}{9} \frac{4!}{4!} - \dots$ (6.9.2) $3 x^7 5 x^{11} x^3$ $S(x) = \frac{p^2}{3} \frac{1!}{1!} - \frac{p^2}{7} \frac{3!}{3!} + \frac{p^2}{11} \frac{5!}{5!} - \dots$ There is a complex continued fraction that yields both $S(x)$ and $C(x)$ simultaneously: $\frac{C(x)+iS(x)}{1+i} = \frac{2}{2} \operatorname{erf} z$, $z = \frac{2}{2} (1-i)x$ (6.9.3) where $1 \frac{1}{2} 1 \frac{3}{2} 2$ $\frac{1}{2} \operatorname{erfc} z = \frac{1}{z} + \frac{z}{z^2} + \frac{z}{z^3} + \dots$

$$z + \frac{1}{2} \frac{1}{z} - \frac{1}{2} \frac{1}{z^3} + \frac{1}{2} \frac{1}{z^5} - \frac{1}{2} \frac{1}{z^7} + \dots \quad (6.9.4)$$

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$$p_{ll} = (x^{2*ll-1})^{*p_{mmp1}} (ll+m-1)^{*p_{mmp1}} / (ll-m)^{*p_{mmp1}}$$

$$p_{mmp1} = p_{mmp1} / p_{mmp1}$$

$$p_{ll} = p_{ll} / p_{ll}$$

END CITED REFERENCES AND FURTHER READING: Magnus, W., and Oberhettinger, F. 1949, *Formulas and Theorems for the Functions of Mathematical Physics* (New York: Chelsea), pp. 54ff. [1] Abramowitz, M., and Stegun, I.A. 1964, *Handbook of Mathematical Functions*, Applied Mathematics Series, Volume 55 (Washington: National Bureau of Standards; reprinted 1968 by Dover Publications, New York), Chapter 8. [2] 6.9 Fresnel Integrals, Cosine and Sine Integrals

Fresnel Integrals The two Fresnel integrals are defined by

$$C(x) = \int_0^x \cos t^2 dt \quad (6.9.1)$$

$$S(x) = \int_0^x \sin t^2 dt$$

The most convenient way of evaluating these functions to arbitrary precision is to use power series for small x and a continued fraction for large x . The series are

$$C(x) = x - \frac{1}{2} \frac{x^5}{5!} + \frac{1}{2} \frac{x^9}{9!} - \dots \quad (6.9.2)$$

$$S(x) = \frac{1}{2} \frac{x^3}{3!} - \frac{1}{2} \frac{x^7}{7!} + \frac{1}{2} \frac{x^{11}}{11!} - \dots$$

There is a complex continued fraction that yields both $S(x)$ and $C(x)$ simultaneously:

$$C(x) + iS(x) = \frac{1+i}{2} \operatorname{erf} z, \quad z = \frac{1-i}{2} x \quad (6.9.3)$$

where

$$\operatorname{erfc} z = \frac{1}{z} + \frac{z}{z^2 + 1} + \frac{z^2}{z^2 + 4} + \frac{z^3}{z^2 + 9} + \dots \quad (6.9.4)$$

6.9 Fresnel Integrals, Cosine and Sine Integrals 249

In the last line we have converted the standard form of the continued fraction to its even form (see 5.2), which converges twice as fast. We must be careful not to evaluate the alternating series (6.9.2) at too large a value of x ; inspection of the terms shows that $x = 1.5$ is a good point to switch over to the continued fraction. Note that for large x

$$C(x) \sim \frac{1}{2} + \frac{1}{2} \frac{\pi x \sin p^2}{2} - \frac{1}{2} \frac{\pi x \cos p^2}{2} x^2, \quad S(x) \sim \frac{1}{2} - \frac{1}{2} \frac{\pi x^2}{2} \quad (6.9.5)$$

Thus the precision of the routine `frenel` may be limited by the precision of the library routines for sine and cosine for large x . SUBROUTINE `frenel(x,s,c)` INTEGER

MAXIT REAL c,s,x,EPS,FPMIN,PI,PIBY2,XMIN PARAMETER (EPS=6.e-8,MAXIT=100,FPMIN=1.e-30,XMIN=1.5, *PI=3.1415927,PIBY2=1.5707963) Computes the Fresnel integrals S(x) and C(x) for all real x. Parameters: EPS is the relative error; MAXIT is the maximum number of iterations allowed; FPMIN is a number near the smallest representable floating-point number; XMIN is the dividing line between using the series and continued fraction; PI = p; PIBY2 = p/2. INTEGER k,n REAL a,absc,ax,act,pix2,sign,sum,sumc,sums,term,test COMPLEX b,cc,d,h,del,cs LOGICAL odd absc(h)=abs(real(h))+abs(aimag(h)) Statement function. ax=abs(x) if(ax.lt.sqrt(FPMIN))then Special case: avoid failure of convergence test because of underflow. s=0. c=ax else if(ax.le.XMIN)then Evaluate both series simultaneously. sum=0. sums=0. sumc=ax sign=1. fact=PIBY2*ax*ax odd=.true. term=ax n=3 do 1 1 k=1,MAXIT term=term*fact/k sum=sum+sign*term/n test=abs(sum)*EPS if(odd)then sign=-sign sums=sum sum=sumc else sumc=sum sum=sums endif if(term.lt.test)goto 1 odd=.not.odd n=n+2 enddo 1 1 pause series failed in fresnel 1 s=sums c=sumc else Evaluate continued fraction by modified Lentz's method (5.2). pix2=PI*ax*ax b=cplx(1.,-pix2)

250 Chapter 6. Special Functions cc=1./FPMIN d=1./b h=d n=-1 do 1 2 k=2,MAXIT n=n+2 a=-n*(n+1) b=b+4. d=1./(a*d+b) Denominators cannot be zero. cc=b+a/cc del=cc*d h=h*del if(absc(del-1.).lt.EPS)goto 2 enddo 1 2 pause cf failed in fresnel 2 h=h*cplx(ax,-ax) cs=cplx(.5,.5)*(1.-cplx(cos(.5*pix2),sin(.5*pix2))*h) c=real(cs) s=aimag(cs) endif if(x.lt.0.)then Use antisymmetry. c=-c s=-s endif return END Cosine and Sine Integrals The cosine and sine integrals are defined by $\cos t - 1$ $Ci(x) = \int_0^x \frac{t}{t} dt$ (6.9.6) $0 \sin t$ $Si(x) = \int_0^x \frac{t}{t} dt$ 0 Here 0.5772... is Euler's constant. We only need a way to calculate the functions for $x \geq 0$, because $Si(-x) = -$

$\text{Si}(x)$, $\text{Ci}(-x) = \text{Ci}(x) - i\pi$ (6.9.7) Once again we can evaluate these functions by a judicious combination of power series and complex continued fraction. The series are $\text{Si}(x) = x - \frac{x^3}{3!} + \frac{x^5}{5!} - \frac{x^7}{7!} + \frac{x^9}{9!} - \dots$ (6.9.8) $\text{Ci}(x) = \gamma + \ln x - \frac{x^2}{4 \cdot 4!} + \frac{x^4}{2 \cdot 2!} - \frac{x^6}{6 \cdot 6!} + \frac{x^8}{4 \cdot 4!} - \dots$ The continued fraction for the exponential integral $E_1(ix)$ is $E_1(ix) = -\text{Ci}(x) + i[\text{Si}(x) - \pi/2]$ (6.9.9) $\frac{1}{ix + 1} = \frac{1}{ix + 1} - \frac{1}{2} \frac{1}{ix + 1} + \frac{1}{2} \frac{1}{ix + 1} - \frac{1}{2} \frac{1}{ix + 1} + \dots$

6.9 Fresnel Integrals, Cosine and Sine Integrals 251 The even form of the continued fraction is given in the last line and converges twice as fast for about the same amount of computation. A good crossover point from the alternating series to the continued fraction is $x = 2$ in this case. As for the Fresnel integrals, for large x the precision may be limited by the precision of the sine and cosine routines.

```

SUBROUTINE cisi(x,ci,si) INTEGER MAXIT REAL ci,si,x,EPS,EULER,PIBY2,FPMIN,TMIN
PARAMETER (EPS=6.e-8,EULER=.57721566,MAXIT=100,PIBY2=1
* FPMIN=1.e-30,TMIN=2.) Computes the cosine and sine
integrals Ci(x) and Si(x). Ci(0) is returned as a large
negative number and no error message is generated. For
x<0 the routine returns Ci(-x) and you must supply the -
ip yourself. Parameters: EPS is the relative error, or ab-
solute error near a zero of Ci(x); EULER = ; MAXIT
is the maximum number of iterations allowed; PIBY2 =
p/2; FPMIN is a number near the smallest representable
floating-point number; TMIN is the dividing line between
using the series and continued fraction. INTEGER i,k
REAL a,err,fact,sign,sum,sumc,sums,t,term,absc COMPLEX
h,b,c,d,del LOGICAL odd absc(h)=abs(real(h))+abs(aimag(h))
Statement function. t=abs(x) if(t.eq.0.)then Special case.
si=0. ci=-1./FPMIN return endif if(t.gt.TMIN)then Eval-
uate continued fraction by modified Lentzs method (5.2).
b=cmplx(1.,t) c=1./FPMIN d=1./b h=d do 1 1 i=2,MAXIT

```

```

a=- (i-1)**2 b=b+2. d=1./(a*d+b) Denominators cannot be
zero. c=b+a/c del=c*d h=h*del if(abs(c-del-1.)<.EPS)goto 1
enddo 1 1 pause cf failed in cisi 1 continue h=cplx(cos(t),-
sin(t))*h ci=-real(h) si=PI/2+aimag(h) else Evaluate both
series simultaneously. if(t<.sqrt(FPMIN))then Special case:
avoid failure of convergence test because of under ow. sumc=0.
sums=t else sum=0. sums=0. sumc=0. sign=1. fact=1.
odd=.true. do 1 2 k=1,MAXIT fact=fact*t/k term=fact/k

```

```

252 Chapter 6. Special Functions sum=sum+sign*term
err=term/abs(sum) if(odd)then sign=-sign sums=sum sum=sumc
else sumc=sum sum=sums endif if(err<.EPS)goto 2 odd=.not.odd

```

```

enddo 1 2 pause maxits exceeded in cisi endif 2 si=sums
ci=sumc+log(t)+EULER endif if(x<.0.)si=-si return END

```

CITED REFERENCES AND FURTHER READING: Stegun,

I.A., and Zucker, R. 1976, Journal of Research of the National Bureau of Standards, vol. 80B, pp. 291311; 1981, op. cit.,

vol. 86, pp. 661686. Abramowitz, M., and Stegun, I.A. 1964, Handbook of Mathematical Functions, Applied Mathematics Series, Volume 55 (Washington: National Bureau of Standards;

reprinted 1968 by Dover Publications, New York), Chapters 5 and 7.

6.10 Dawson's Integral Dawson's Integral $F(x)$ is defined by $x \int_0^x e^{-t^2} dt$ (6.10.1) $F(x) = e^{-x^2} \int_0^x e^{t^2} dt$ The function can also be related to the complex error function by $F(z) = \frac{1}{2} e^{-z^2} [1 - \operatorname{erfc}(-iz)]$.

(6.10.2) A remarkable approximation for $F(z)$, due to Rybicki [1], is $F(z) = \lim_{n \rightarrow \infty} e^{-(z-nh)^2} \sum_{k=0}^{n-1} \frac{1}{k!} (zh)^k$ (6.10.3)

$h > 0$ What makes equation (6.10.3) unusual is that its accuracy increases exponentially as h gets small, so that quite moderate values of h (and correspondingly quite rapid convergence of the series) give very accurate approximations.

252 Chapter 6. Special Functions $sum = sum + sign * term$
 $err = term / abs(sum)$ if(odd)then $sign = -sign$ $sums = sum$ $sum = sumc$
else $sumc = sum$ $sum = sums$ endif if($err < .EPS$)goto 2 odd=.not.odd

252 Chapter 6. Special Functions $sum = sum + sign * term$
 $err = term / abs(sum)$ if(odd)then $sign = -sign$ $sums = sum$ $sum = sumc$
else $sumc = sum$ $sum = sums$ endif if($err < .EPS$)goto 2 odd=.not.odd

enddo 1 2 pause maxits exceeded in cisi endif 2 si=sums
 ci=sumc+log(t)+EULER endif if(x.lt.0.)si=-si return END
 CITED REFERENCES AND FURTHER READING: Stegun,
 I.A., and Zucker, R. 1976, Journal of Research of the National
 Bureau of Standards, vol. 80B, pp. 291311; 1981, op. cit.,
 vol. 86, pp. 661686. Abramowitz, M., and Stegun, I.A. 1964,
 Handbook of Mathematical Functions, Applied Mathe-
 matics Series, Volume 55 (Washington: National Bureau of Standards;
 reprinted 1968 by Dover Publications, New York), Chapters 5
 and 7. 6.10 Dawsons Integral Dawsons Integral $F(x)$ is defined
 by $x \int_0^\infty e^{-x^2 t^2} dt$ (6.10.1) $F(x) = e^{-x^2} \int_0^\infty e^{-t^2} dt$ The function can also be re-
 lated to the complex error function by $F(z) = \frac{1}{2} e^{-z^2} [1 - \text{erfc}(-iz)]$. (6.10.2) A remarkable approximation for $F(z)$, due
 to Rybicki [1], is $F(z) = \lim_{h \rightarrow 0} e^{-z^2} \sum_{n=0}^{\infty} \frac{(-1)^n (z/h)^{2n}}{(2n)!}$ (6.10.3)
 $h \rightarrow 0$ What makes equation (6.10.3) unusual is that its accuracy
 increases exponentially as h gets small, so that quite moderate
 values of h (and correspondingly quite rapid convergence of the
 series) give very accurate approximations.

253 6.10 Dawsons Integral We will discuss the theory that
 leads to equation (6.10.3) later, in 13.11, as an interesting appli-
 cation of Fourier methods. Here we simply implement a routine
 for real values of x based on the formula. It is first convenient
 to shift the summation index to center it approximately on the
 maximum of the exponential term. Define n_0 to be the even
 integer nearest to x/h , and $x_0 = n_0 h$, $x = x - x_0$, and $n = n$
 $- n_0$, so that $N F(x) = \sum_{n=-N}^N e^{-(x-nh)^2}$, (6.10.4) $n + n_0$
 N where the approximate equality is accurate when h is
 sufficiently small and N is sufficiently large. The computation
 of this formula can be greatly speeded up if we note that $e^{-(x-nh)^2} = e^{-x^2} e^{-(nh)^2} e^{2xnh}$. (6.10.5) The first factor
 is computed once, the second is an array of constants to be
 stored, and the third can be computed recursively, so that only

two exponentials need be evaluated. Advantage is also taken of the symmetry of the coefficients $e^{-(nh)^2}$ by breaking the summation up into positive and negative values of n separately. In the following routine, the choices $h = 0.4$ and $N = 11$ are made. Because of the symmetry of the summations and the restriction to odd values of n , the limits on the do loops are 1 to 6. The accuracy of the result in this REAL version is about 2×10^{-7} . In order to maintain relative accuracy near $x = 0$, where $F(x)$ vanishes, the program branches to the evaluation of the power series [2] for $F(x)$, for $|x| < 0.2$.

```

FUNCTION dawson(x) IN-
TEGGER NMAX REAL dawson,x,H,A1,A2,A3 PARAMETER
(NMAX=6,H=0.4,A1=2./3.,A2=0.4,A3=2./7.) Returns Daw-
sons integral  $F(x) = \int_0^x \exp(-t^2) dt$  for any real x.
INTEGGER i,init,n0 REAL d1,d2,e1,e2,sum,x2,xp,xx,c(NMAX)
SAVE init,c DATA init/0/ Flag is 0 if we need to initialize,
else 1. if(init.eq.0)then init=1 do 1 1 i=1,NMAX c(i)=exp(-
((2.*float(i)-1.)*H)**2) enddo 1 1 endif if(abs(x).lt.0.2)then
Use series expansion. x2=x**2 dawson=x*(1.-A1*x2*(1.-
A2*x2*(1.-A3*x2))) else Use sampling theorem representa-
tion. xx=abs(x) n0=2*nint(0.5*xx/H) xp=xx-float(n0)*H
e1=exp(2.*xp*H) e2=e1**2 d1=float(n0+1) d2=d1-2. sum=0.
do 1 2 i=1,NMAX

```

254 Chapter 6. Special Functions $sum = sum + c(i) * (e1/d1 + 1./((d2*e1)$
 $d1 = d1 + 2. d2 = d2 - 2. e1 = e2 * e1$ enddo 1 2 $dawson = 0.5641895835 * sign(\exp$
 $xp**2), x) * sum$ Constant is $1/vp$. endif return END Other
methods for computing Dawsons integral are also known [2,3].
CITED REFERENCES AND FURTHER READING: Rybicki,
G.B. 1989, Computers in Physics, vol. 3, no. 2, pp. 85-87. [1]
Cody, W.J., Pociorek, K.A., and Thatcher, H.C. 1970, Mathe-
matics of Computation, vol. 24, pp. 171-178. [2] McCabe, J.H.
1974, Mathematics of Computation, vol. 28, pp. 811-816. [3]
6.11 Elliptic Integrals and Jacobian Elliptic Functions Elliptic

integrals occur in many applications, because any integral of the form $\int R(t, s) dt$ (6.11.1) where R is a rational function of t and s , and s is the square root of a cubic or quartic polynomial in t , can be evaluated in terms of elliptic integrals. Standard references [1] describe how to carry out the reduction, which was originally done by Legendre. Legendre showed that only three basic elliptic integrals are required. The simplest of these is $\int dt / \sqrt{(a_1 + b_1 t)(a_2 + b_2 t)(a_3 + b_3 t)(a_4 + b_4 t)}$ (6.11.2) where we have written the quartic s^2 in factored form. In standard integral tables [2], one of the limits of integration is always a zero of the quartic, while the other limit lies closer than the next zero, so that there is no singularity within the interval. To evaluate I_1 , we simply break the interval $[y, x]$ into subintervals, each of which either begins or ends on a singularity. The tables, therefore, need only distinguish the eight cases in which each of the four zeros (ordered according to size) appears as the upper or lower limit of integration. In addition, when one of the b s in (6.11.2) tends to zero, the quartic reduces to a cubic, with the largest or smallest singularity moving to ∞ ; this leads to eight more cases (actually just special cases of the first eight). The sixteen cases in total are then usually tabulated in terms of Legendre's standard elliptic integral of the 1st kind, which we will define below. By a change of the variable of integration t , the zeros of the quartic are mapped to standard locations

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```

sum=sum+c(i)*(e1/d1+1./(d2*e1)
d1=d1+2. d2=d2-2. e1=e2*e1 enddo 1 2
dawson=0.5641895835*sign(exp
xp**2),x)*sum
Constant is 1/vp.
endif return END

```

Other methods for computing Dawson's integral are also known [2,3].
CITED REFERENCES AND FURTHER READING: Rybicki, G.B. 1989, *Computers in Physics*, vol. 3, no. 2, pp. 85-87. [1] Cody, W.J., Pociorek, K.A., and Thatcher, H.C. 1970, *Mathe-*

matics of Computation, vol. 24, pp. 171178. [2] McCabe, J.H. 1974, Mathematics of Computation, vol. 28, pp. 811816. [3]

6.11 Elliptic Integrals and Jacobian Elliptic Functions

Elliptic integrals occur in many applications, because any integral of the form $\int R(t, s) dt$ (6.11.1) where R is a rational function of t and s , and s is the square root of a cubic or quartic polynomial in t , can be evaluated in terms of elliptic integrals. Standard references [1] describe how to carry out the reduction, which was originally done by Legendre. Legendre showed that only three basic elliptic integrals are required. The simplest of these is $\int dt / \sqrt{(a_1 + b_1 t)(a_2 + b_2 t)(a_3 + b_3 t)(a_4 + b_4 t)}$ (6.11.2) where we have written the quartic s^2 in factored form. In standard integral tables [2], one of the limits of integration is always a zero of the quartic, while the other limit lies closer than the next zero, so that there is no singularity within the interval. To evaluate I_1 , we simply break the interval $[y, x]$ into subintervals, each of which either begins or ends on a singularity. The tables, therefore, need only distinguish the eight cases in which each of the four zeros (ordered according to size) appears as the upper or lower limit of integration. In addition, when one of the b s in (6.11.2) tends to zero, the quartic reduces to a cubic, with the largest or smallest singularity moving to ∞ ; this leads to eight more cases (actually just special cases of the first eight). The sixteen cases in total are then usually tabulated in terms of Legendre's standard elliptic integral of the 1st kind, which we will define below. By a change of the variable of integration t , the zeros of the quartic are mapped to standard locations

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 on the real axis. Then only two dimensionless parameters are needed to tabulate Legendre's integral. However, the symmetry of the original integral (6.11.2) under permutation of the roots

is concealed in Legendres notation. We will get back to Legendres notation below. But first, here is a better way: Carlson [3] has given a new definition of a standard elliptic integral of the first kind,
$$F(x, y, z) = \int_0^1 \frac{dt}{\sqrt{(t+x)(t+y)(t+z)}} \quad (6.11.3)$$
 where $x, y,$ and z are nonnegative and at most one is zero. By standardizing the range of integration, he retains permutation symmetry for the zeros. (Weierstrass canonical form also has this property.) Carlson first shows that when x or y is a zero of the quartic in (6.11.2), the integral F can be written in terms of R_F in a form that is symmetric under permutation of the remaining three zeros. In the general case when neither x nor y is a zero, two such R_F functions can be combined into a single one by an addition theorem, leading to the fundamental formula
$$F(x, y, z) = 2R_F(U_1, U_2, U_3, U_4) \quad (6.11.4)$$
 where $U_i = \sqrt{(x_i - x)(x_j - y)(x_k - z)(x_m - z)}/(x - y)$
$$x_i = (a_i + b_i x)^{1/2}, y_i = (a_i + b_i y)^{1/2} \quad (6.11.5)$$
 and i, j, k, m is any permutation of $1, 2, 3, 4$. A short-cut in evaluating these expressions is
$$U_1^2 U_2^2 - (a_1 b_4 - a_4 b_1)(a_2 b_3 - a_3 b_2) = U_3^2 \quad (6.11.6)$$

$$U_1^2 U_2^2 - (a_1 b_3 - a_3 b_1)(a_2 b_4 - a_4 b_2) = U_4^2$$
 The U_i correspond to the three ways of pairing the four zeros, and F is thus manifestly symmetric under permutation of the zeros. Equation (6.11.4) therefore reproduces all sixteen cases when one limit is a zero, and also includes the cases when neither limit is a zero. Thus Carlsons function allows arbitrary ranges of integration and arbitrary positions of the branch points of the integrand relative to the interval of integration. To handle elliptic integrals of the second and third kind, Carlson defines the standard integral of the third kind as
$$R_J(x, y, z, p) = \int_0^1 \frac{dt}{(t+p)\sqrt{(t+x)(t+y)(t+z)}} \quad (6.11.7)$$
 which is symmetric in $x, y,$ and z . The degenerate case when two arguments are equal is denoted $R_D(x, y, z) = R_J(x, y, z, z)$
$$(6.11.8)$$
 and is symmetric in x and y . The function R_D replaces Legendres integral

of the second kind. The degenerate form of RF is denoted RC
 $(x, y) = \text{RF}(x, y, y)$ (6.11.10) It embraces logarithmic, inverse
circular, and inverse hyperbolic functions. Carlson [4-7] gives
integral tables in terms of the exponents of the linear factors
of the quartic in (6.11.1). For example, the integral where the
exponents are $(1/2, -1/2, -3/2, 2)$, 1 can be expressed as a single
integral in terms of RD ; it accounts for 144 separate cases in
Gradshteyn and Ryzhik [2]! Refer to Carlsons papers [3-7] for
some of the practical details in reducing elliptic integrals to his
standard forms, such as handling complex conjugate zeros.

256 Chapter 6. Special Functions Turn now to the numeri-
cal evaluation of elliptic integrals. The traditional methods [8]
are Gauss or Landen transformations. Descending transforma-
tions decrease the modulus k of the Legendre integrals towards
zero, increasing transformations increase it towards unity. In
these limits the functions have simple analytic expressions.
While these methods converge quadratically and are quite sat-
isfactory for integrals of the first and second kinds, they gener-
ally lead to loss of significant figures in certain regimes for inte-
grals of the third kind. Carlsons algorithms [9,10], by contrast,
provide a unified method for all three kinds with no significant
cancellations. The key ingredient in these algorithms is the du-
plication theorem: $\text{RF}(x, y, z) = 2\text{RF}(x + \sqrt{xy}, y + \sqrt{xy}, z + \sqrt{yz})$ (6.11.11)
 $x + \sqrt{xy} = \text{RF}(4x, y + 4x, z + 4x)$ where $\sqrt{xy} = (xy)^{1/2} + (xz)^{1/2} + (yz)^{1/2}$
(6.11.12) This theorem can be proved by a simple change of
variable of integration[11]. Equation (6.11.11) is iterated until
the arguments of RF are nearly equal. For equal arguments
we have $\text{RF}(x, x, x) = x^{-1/2}$ (6.11.13) When the arguments
are close enough, the function is evaluated from a fixed Taylor
expansion about (6.11.13) through fifth-order terms. While the
iterative part of the algorithm is only linearly convergent, the
error ultimately decreases by a factor of $46 = 4096$ for each

iteration. Typically only two or three iterations are required, perhaps six or seven if the initial values of the arguments have huge ratios. We list the algorithm for RF here, and refer you to Carlsons paper [9] for the other cases. Stage 1: For $n = 0, 1, 2, \dots$ compute $x_n = (x_n + y_n + z_n) / 3$, $Y_n = 1 - (x_n / n)$, $Z_n = 1 - (y_n / n)$, $Z_n = 1 - (z_n / n)$, $n = \max(-X_n, -Y_n, -Z_n)$. If $n \geq \text{tol}$ go to Stage 2; else compute $x_{n+1} = (x_n + n) / 4$, $y_{n+1} = (y_n + n) / 4$, $z_{n+1} = (z_n + n) / 4$ and repeat this stage. Stage 2: Compute $E2 = X_n Y_n - Z_n^2$, $E3 = X_n Y_n Z_n$, $\text{RF} = (1 - 1/2 E2 + 1/4 E3 + 1/2 E2 E3) / (n + 1/2)$. In some applications the argument p in RJ or the argument y in RC is negative, and the Cauchy principal value of the integral is required. This is easily handled by using the formulas $\text{RJ}(x, y, z, p) = [(-y)\text{RJ}(x, y, z, -p) - 3\text{RF}(x, y, z) + 3\text{RC}(xz/y, p/y)] / (y - p)$ (6.11.14) where $p = y + (z - y)(y - x)$ (6.11.15) $y - p$

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 is positive if p is negative, and $1/2 \text{RC}(x - y, -y)$ (6.11.16)

$\text{RC}(x, y) = x \sqrt{x - y}$ The Cauchy principal value of RJ has a zero at some value of $p \neq 0$, so (6.11.14) will give some loss of significant figures near the zero. FUNCTION rf(x,y,z) REAL rf,x,y,z,ERRTOL,TINY,BIG,THIRD,C1,C2,C3,C4 PARAMETER (ERRTOL=.08,TINY=1.5e-38,BIG=3.E37,THIRD=1./3.,

* C1=1./24.,C2=.1,C3=3./44.,C4=1./14.) Computes Carlsons elliptic integral of the first kind, RF(x, y,z). x, y, and z must be nonnegative, and at most one can be zero. TINY must be at least 5 times the machine under ow limit, BIG at most one fifth the machine over ow limit. REAL alamb,ave,delx,dely,delz,e2,e3,sqrtx,sqrty,sqrtz,xt,yt,zt if(min(x,y,z).lt.0. * max(x,y,z).gt.BIG)pause invalid arguments in rf xt=x yt=y zt=z 1 continue sqrtx=sqrt(xt) sqrty=sqrt(yt) sqrtz=sqrt(zt) alamb=sqrtx*(sqrty+sqrtz)+sqrty*sqrtz xt=.25*(xt+alamb) yt=.25*(yt

```

zt=.25*(zt+alamb) ave=THIRD*(xt+yt+zt) delx=(ave-xt)/ave
dely=(ave-yt)/ave delz=(ave-zt)/ave if(max(abs(delx),abs(dely),abs(delz)
1 e2=delx*dely-delz**2 e3=delx*dely*delz rf=(1.+(C1*e2-C2-
C3*e3)*e2+C4*e3)/sqrt(ave) return END

```

A value of 0.08 for the error tolerance parameter is adequate for single precision (7 n , we see that 0.0025 will yield double precision significant digits). Since the error scales as 6 (16 significant digits) and require at most two or three more iterations. Since the coefficients of the sixth-order truncation error are different for the other elliptic functions, these values for the error tolerance should be changed to 0.04 and 0.0012 in the algorithm for RC , and 0.05 and 0.0015 for RJ and RD . As well as being an algorithm in its own right for certain combinations of elementary functions, the algorithm for RC is used repeatedly in the computation of RJ . The Fortran implementations test the input arguments against two machine-dependent constants, TINY and BIG, to ensure that there will be no under ow or over ow during the computation. We have chosen conservative values, corresponding to a machine minimum of $3 \cdot 10^{-39}$ and a machine maximum of $1.7 \cdot 10^38$. You can always extend the range of admissible argument values by using the homogeneity relations (6.11.22), below.

```

FUNCTION rd(x,y,z)
REAL rd,x,y,z,ERRTOL,TINY,BIG,C1,C2,C3,C4,C5,C6 PA-
RAMETER (ERRTOL=.05,TINY=1.e-25,BIG=4.5E21,C1=3./14.,C2=
* C3=9./22.,C4=3./26.,C5=.25*C3,C6=1.5*C4) Computes Carl-
sons elliptic integral of the second kind, RD (x, y, z). x and y
must be nonnegative, and at most one can be zero. z must be
positive. TINY must be at least twice the negative 2/3 power of
the machine over ow limit. BIG must be at most 0.1 ERRTOL
times the negative 2/3 power of the machine under ow limit.
REAL alamb,ave,delx,dely,delz,ea,eb,ec,ed,ee,fac,sqrtx,sqrty,

```

```

if(min(x,y).lt.0..or.min(x+y,z).lt.TINY.or. * max(x,y,z).gt.BIG)pause
invalid arguments in rd xt=x yt=y zt=z sum=0. fac=1. 1 con-
tinue sqrtx=sqrt(xt) sqrty=sqrt(yt) sqrtz=sqrt(zt) alamb=sqrtx*(sqrty+
sum=sum+fac/(sqrtz*(zt+alamb)) fac=.25*fac xt=.25*(xt+alamb)
yt=.25*(yt+alamb) zt=.25*(zt+alamb) ave=.2*(xt+yt+3.*zt)
delx=(ave-xt)/ave dely=(ave-yt)/ave delz=(ave-zt)/ave if(max(abs(delx)
1 ea=delx*dely eb=delz*delz ec=ea-eb ed=ea-6.*eb ee=ed+ec+ec
rd=3.*sum+fac*(1.+ed*(-C1+C5*ed-C6*delz*ee) * +delz*(C2*ee+delz*
C3*ec+delz*C4*ea)))/(ave*sqrt(ave)) return END FUNCTION
rj(x,y,z,p) REAL rj,p,x,y,z,ERRTOL,TINY,BIG,C1,C2,C3,C4,C5,C6,C7,
PARAMETER (ERRTOL=.05,TINY=2.5e-13,BIG=9.E11,C1=3./14.,C
* C3=3./22.,C4=3./26.,C5=.75*C3,C6=1.5*C4,C7=.5*C2,C8=C3+C3)
C USES rc,rf Computes Carlsons elliptic integral of the third
kind, RJ (x,y, z, p). x, y,andz must be nonnegative, and
at most one can be zero. p must be nonzero. If p<0, the
Cauchy principal value is returned. TINY must be at least
twice the cube root of the machine under ow limit, BIG at
most one fifth the cube root of the machine over ow limit.
REAL a,alamb,alpha,ave,b,beta,delp,delx,dely,delz,ea,eb,ec, *
ed,ee,fac,pt,rcx,rho,sqrtx,sqrty,sqrtz,sum,tau,xt, * yt,zt,rc,rf if(min(x,y,z)
* max(x,y,z,abs(p)).gt.BIG)pause invalid arguments in rj sum=0.
fac=1. if(p.gt.0.)then xt=x yt=y zt=z pt=p else xt=min(x,y,z)
zt=max(x,y,z) yt=x+y+z-xt-zt a=1./(yt-p) b=a*(zt-yt)*(yt-
xt) pt=yt+b rho=xt*zt/yt

```

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6.12 Hypergeometric Functions As was discussed in 5.14, a fast, general routine for the the complex hypergeometric function ${}_2F_1(a, b, c; z)$, is difficult or impossible. The function is defined as the analytic continuation of the hypergeometric series, ${}_2F_1(a, b, c; z) = 1 + \frac{ab}{c}z + \frac{a(a+1)b(b+1)}{c(c+1)2!}z^2 + \frac{a(a+1)\dots(a+j-1)b(b+1)\dots(b+j-1)}{c(c+1)\dots(c+j-1)j!}z^j + \dots$ (6.12.1) This series converges only within the unit circle $|z| < 1$ (see [1]), but ones interest in the function is not confined to this region. Section 5.14 discussed the method of evaluating this function by direct path integration in the complex plane. We here merely list the routines that result. Implementation of the function hypgeo is straightforward, and is described by comments in the program. The machinery associated with Chapter 16s routine for integrating differential equations, odeint, is only minimally intrusive, and need not even be completely understood: use of odeint requires a common block with one zeroed variable, one subroutine call, and a

prescribed format for the derivative routine hypdrv. The sub-routine hypgeo will fail, of course, for values of z too close to the singularity at 1. (If you need to approach this singularity, or the one at 8, use the linear transformation formulas in 15.3 of [1].) Away from $z = 1$, and for moderate values of a , b , c , it is often remarkable how few steps are required to integrate the equations. A half-dozen is typical.

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```

FUNCTION hypgeo(a,b,c,z) COMPLEX hypgeo,a,b,c,z REAL EPS PARAMETER
(EPS=1.e-6) Accuracy parameter. C USES bsstep,hypdrv,hypser,odeint
Complex hypergeometric function  ${}_2F_1$  for complex  $a$ ,  $b$ ,  $c$ , and  $z$ ,
by direct integration of the hypergeometric equation in the
complex plane. The branch cut is taken to lie along the
real axis,  $\text{Re } z \geq 1$ . INTEGER kmax,nbad,nok EXTERNAL
bsstep,hypdrv COMPLEX z0,dz,aa,bb,cc,y(2) COMMON /hypg/
aa,bb,cc,z0,dz COMMON /path/ kmax Used by odeint. kmax=0
if (real(z)**2+aimag(z)**2.le.0.25) then Use series... call
hypser(a,b,c,z,hypgeo,y(2)) return else if (real(z).lt.0.) then
...or pick a starting point for the path inte- gration. z0=cmplx(-
0.5,0.) else if (real(z).le.1.0) then z0=cmplx(0.5,0.) else
z0=cmplx(0.,sign(0.5,aimag(z))) endif aa=a Load the common
block, used to pass pa- rameters over the head of odeint to hyp-
drv. bb=b cc=c dz=z-z0 call hypser(aa,bb,cc,z0,y(1),y(2)) Get
starting function and derivative. call odeint(y,4,0.,1.,EPS,.1,.0001,nok,nb)
The arguments to odeint are the vector of independent variabl
es, its length, the starting and ending values of the dependent
variable, the accuracy parameter, an initial guess for stepsize,
a minimum stepsize, the (returned) number of good and bad
steps taken, and the names of the derivative routine and the
(here Bulirsch-Stoer) stepping routine. hypgeo=y(1) return
END SUBROUTINE hypser(a,b,c,z,series,deriv) INTEGER n
COMPLEX a,b,c,z,series,deriv,aa,bb,cc,fac,temp Returns the

```

hypergeometric series ${}_2F_1$ and its derivative, iterating to machine accuracy. For $\text{cabs}(z) = 1/2$ convergence is quite rapid.

```

deriv=cplx(0.,0.) fac=cplx(1.,0.) temp=fac aa=a bb=b
cc=c do 1 1 n=1,1000 fac=((aa*bb)/cc)*fac deriv=deriv+fac
fac=fac*z/n series=temp+fac if (series.eq.temp) return temp=series
aa=aa+1. bb=bb+1. cc=cc+1. enddo 1 1 pause convergence
failure in hypser END

```

265 6.12 Hypergeometric Functions SUBROUTINE hyp-
 drv(s,y,dyds) REAL s COMPLEX y(2),dyds(2),aa,bb,cc,z0,dz,z
 Derivative subroutine for the hypergeometric equation, see
 text equation (5.14.4). COMMON /hypg/ aa,bb,cc,z0,dz
 $z=z0+s*dz$ $dyds(1)=y(2)*dz$ $dyds(2)=((aa*bb)*y(1)-(cc-((aa+bb)+1.)*z))$
 return END CITED REFERENCES AND FURTHER
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 ries, Volume 55 (Washington: National Bureau of Standards;
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